# One-dimensional flow of a compressible viscous micropolar fluid: The Cauchy problem

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**Abstract**. The Cauchy problem for one-dimensional flow of a compressible viscous heat-conducting micropolar fluid is considered. It is assumed that the fluid is thermodynamically perfect and polytropic. A corresponding initial-boundary value problem has a unique strong solution on  $]0,1[\times]0,T[$ , for each T > 0. By using this result we construct a sequence of approximate solutions which converges to a solution of the Cauchy problem.

**Key words:** *micropolar fluid, the Cauchy problem, strong solution, weak convergence* 

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## 1. Statement of the problem and the main result

In this paper we consider nonstationary 1-D flow of a compressible and heatconducting micropolar fluid. The equations of motion for this fluid are derived from the integral form of conservation laws for polar fluids, under a number of supplementary assumptions such as politropy, Fourier's law, Boyle's law and selection of constitutive equations (see [7]). A corresponding initial-boundary value problem has a unique strong solution on  $]0,1[\times]0,T[$ , for each T > 0 ([8]). By using this result we prove a global-in-time existence theorem for the Cauchy problem. In our proof we follow some ideas of S.N.Antontsev, A.V.Kazhykhov and V.N.Monakhov, applied to the case of a classical fluid ([1]).

Let  $\rho, v, \omega$  and  $\theta$  denote, respectively, the mass density, velocity, microrotation velocity and temperature of the fluid in the Lagrangean description. Governing equations of the flow under consideration are as follows ([7]):

$$\frac{\partial \rho}{\partial t} + \rho^2 \frac{\partial v}{\partial x} = 0 , \qquad (1.1)$$

$$\frac{\partial v}{\partial t} = \frac{\partial}{\partial x} \left( \rho \frac{\partial v}{\partial x} \right) - K \frac{\partial}{\partial x} \left( \rho \theta \right) , \qquad (1.2)$$

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$$\rho \frac{\partial \omega}{\partial t} = A \left[ \rho \frac{\partial}{\partial x} \left( \rho \frac{\partial \omega}{\partial x} \right) - \omega \right] , \qquad (1.3)$$

$$\rho \frac{\partial \theta}{\partial t} = -K\rho^2 \theta \frac{\partial v}{\partial x} + \rho^2 \left(\frac{\partial v}{\partial x}\right)^2 + \rho^2 \left(\frac{\partial \omega}{\partial x}\right)^2 + \omega^2 + D\rho \frac{\partial}{\partial x} \left(\rho \frac{\partial \theta}{\partial x}\right)$$
(1.4)

in  $\mathbb{R} \times \mathbb{R}^+$ , where K, A and D are positive constants. Equations (1.1)-(1.4) are, respectively, local forms of the conservation laws for the mass, momentum, momentum moment and energy. We take the following non-homogeneous initial conditions:

$$\rho(x,0) = \rho_0(x) , \qquad (1.5)$$

$$v(x,0) = v_0(x) , (1.6)$$

$$\omega(x,0) = \omega_0(x) , \qquad (1.7)$$

$$\theta(x,0) = \theta_0(x) \tag{1.8}$$

for  $x \in \mathbb{R}$ , where  $\rho_0, v_0, \omega_0$  and  $\theta_0$  are given functions. We assume that there exist  $m, M \in \mathbb{R}^+$ , such that

$$0 < m \le \rho_0(x) \le M, \quad m \le \theta_0(x) \le M, \quad x \in \mathbb{R}.$$
(1.9)

The aim of this paper is to prove the following theorem.

**Theorem 1.1.** Let the initial functions satisfy conditions (1.9) and

$$\rho_0 - 1, v_0, \omega_0, \theta_0 - 1 \in H^1(R).$$
(1.10)

Then for each  $T \in \mathbb{R}^+$  there exists a state function

$$S(x,t) = (\rho, v, \omega, \theta)(x,t) \quad (x,t) \in \Pi = \mathbb{R} \times ]0, T[, \qquad (1.11)$$

with the properties:

$$\rho - 1 \in L^{\infty}(0, T; H^1(R)) \cap H^1(\Pi) , \qquad (1.12)$$

$$v, \omega, \theta - 1 \in L^{\infty}(0, T; H^1(R)) \cap H^1(\Pi) \cap L^2(0, T; H^2(R))$$
 (1.13)

which satisfies equations (1.1)-(1.4) in the sense of distributions in  $\Pi$  and conditions (1.5)-(1.8) in the sense of traces.

We denote by  $B^k(R), k \in \mathbf{N}_0$ , the Banach space

$$B^{k}(R) = \{ u \in C^{k}(R) : \lim_{|x| \to \infty} |D^{n}u(x)| = 0, \ 0 \le n \le k \},$$
(1.14)

where  $D^n$  is the *n*-th derivative; the norm is defined by

$$||u||_{B^k(R)} = \sup_{n \le k} \{ \sup_{x \in R} |D^n u(x)| \}.$$
(1.15)

**Remark 1.1.** From Sobolev's embedding theorem ([3, Chapter IV]) and theory of vector-valued distributions ([4, pp. 467-480]) one can conclude that from (1.12) and (1.13) it follows:

$$\rho - 1 \in L^{\infty}(0, T; B^0(R)) \cap C([0, T]; L^2(R)),$$
(1.16)

$$v, \omega, \theta - 1 \in L^2(0, T; B^1(R)) \cap C([0, T]; H^1(R)) \cap L^\infty(0, T; B^0(R))$$
(1.17)

and hence

$$v, \omega \in C([0,T]; B^0(R)), \quad \rho, \theta \in L^{\infty}(\Pi).$$

$$(1.18)$$

The state function S and its distributional derivatives that occur in (1.1)-(1.4) are locally integrable functions in  $\Pi$  and system (1.1)-(1.4) is satisfied a. e. in  $\Pi$ . In other words, state function (1.11) is a strong solution of our system (1.1)-(1.4).

In the proof of *Theorem 1.1* we construct a sequence of approximations  $S_n = (\rho_n, v_n, \omega_n, \theta_n)_{n \in \mathbb{N}}$  of the state function in  $\Pi$ ; we establish some estimates of approximations  $S_n$  which show that  $\{S_n\}_{n \in \mathbb{N}}$  belongs to a fixed ball (i.e. independent of n) of a certain normed space. Using the results of weak compactness of a unit ball in a Hilbert space (resp. a Banach space or resp. the dual of a normed space) from  $\{S_n\}_{n \in \mathbb{N}}$  we extract a subsequence which has limit in the some weak sense. Finally, we show that this limit is the solution of our problem.

## 2. Approximate solutions and a priori estimates

First we introduce the restrictions of the initial functions  $\rho_0$  and  $\theta_0$  to ]-n, n[. For  $n \in \mathbf{N}$  let

$$\rho_{0n} = \rho_0 \text{ on } ]-n, n[,$$
(2.1)

$$\theta_{0n} = \theta_0 \quad \text{on} \quad ]-n, n[. \tag{2.2}$$

We can easily verify that

$$\rho_{0n}, \ \theta_{0n} \in H^1(]-n, n[).$$
(2.3)

Since  $\mathcal{D}(R)$  is dense in  $H^1(R)$ , there exist the sequences  $\{v_{0n}\}$  and  $\{\omega_{0n}\}$  of approximations of the initial functions  $v_0$  and  $\omega_0$  with the following properties:

(i) 
$$v_{0n}, \omega_{0n} \in H^1_0(] - n, n[), \ v_{0n} = 0, \ \omega_{0n} = 0 \text{ on } \mathbb{R} \setminus ] - n, n[, (2.4)$$

(ii) 
$$v_{0n} \to v_0, \ \omega_{0n} \to \omega_0 \text{ strongly in } H^1(R).$$
 (2.5)

Let us consider equations (1.1) - (1.4) on  $] - n, n[\times \mathbb{R}^+$ , with the boundary conditions

$$v(-n,t) = v(n,t) = 0, \ \omega(-n,t) = \omega(n,t) = 0,$$
 (2.6)

$$\frac{\partial\theta}{\partial x}(-n,t) = \frac{\partial\theta}{\partial x}(n,t) = 0$$
(2.7)

for t > 0 and with  $(\rho_{0n}, v_{0n}, \omega_{0n}, \theta_{0n})$  as the initial data on ] - n, n[. Functions  $\rho_{0n}, v_{0n}, \omega_{0n}$  and  $\theta_{0n}$  satisfy the conditions taken for the initial functions in Theorem 1.1 of [8] and we conclude that, for each  $n \in \mathbf{N}$  and T > 0, problem (1.1)-(1.4), (2.6)-(2.7) has a unique strong solution

$$S_n(x,t) = (\rho_n, v_n, \omega_n, \theta_n)(x,t) , \ (x,t) \in Q_{nT} = ] - n, n[\times]0, T[$$
(2.8)

with the properties:

$$\rho_n \in L^{\infty}(0,T; H^1(]-n, n[)) \cap H^1(Q_{nT}) , \qquad (2.9)$$

$$v_n, \omega_n, \theta_n \in L^{\infty}(0, T; H^1(] - n, n[)) \cap H^1(Q_{nT})$$
  
 
$$\cap L^2(0, T; H^2(] - n, n[)) , \qquad (2.10)$$

$$\rho_n > 0 , \ \theta_n > 0 \quad \text{on } \bar{Q}_{nT}, \tag{2.11}$$

that in [7] and [8] is named a generalised solution. From embedding theorem ([6, Chapter II, Theorem 2.2.1]), theories of vector-valued distributions and interpolations ([4, pp. 467-480]) we observe that from (2.9) and (2.10) it follows:

$$\rho_n \in L^{\infty}(0,T; C([-n,n])) \cap C([0,T]; L^2(]-n,n[),$$
(2.12)

$$v_n, \omega_n, \theta_n \in L^2(0, T; C^1([-n, n])) \cap C([0, T]; H^1(] - n, n[)),$$
 (2.13)

$$v_n, \omega_n, \theta_n \in C(\bar{Q}_{nT}). \tag{2.14}$$

From the properties of the function  $\rho_n$  (see [1, pp. 44-45]) we get

$$\rho_n \in C(\bar{Q}_{nT}). \tag{2.15}$$

Next we prove uniform (in  $n \in \mathbf{N}$ ) a priori estimates for  $S_n$  in  $Q_{nT}$ . By  $C \in \mathbb{R}^+$  we denote a generic constant, independent of  $n \in \mathbf{N}$ .

We introduce non-negative functions  $U_n$  and  $V_n$  defined on ]0, T[ by

$$U_n(t) = \int_{-n}^{n} \left[ \frac{1}{2K} v_n^2 + \frac{1}{2AK} \omega_n^2 + \frac{1}{\rho_n} (\rho_n \ln \rho_n - \rho_n + 1) + \frac{1}{K} (\theta_n - \ln \theta_n - 1) \right] dx, \qquad (2.16)$$

$$V_n(t) = \frac{1}{K} \int_{-n}^{n} \left[ \frac{\rho_n}{\theta_n} \left( \frac{\partial v_n}{\partial x} \right)^2 + \frac{\rho_n}{\theta_n} \left( \frac{\partial \omega_n}{\partial x} \right)^2 + \frac{\omega_n^2}{\rho_n \theta_n} + D \frac{\rho_n}{\theta_n^2} \left( \frac{\partial \theta_n}{\partial x} \right)^2 \right] dx.$$
(2.17)

Using the inequality  $\ln x \leq x - 1$  for  $U_n(0)$  we have

$$U_n(0) \le \int_R \left[ \frac{1}{2K} v_{0n}^2 + \frac{1}{2AK} \omega_{0n}^2 + \frac{(\rho_0 - 1)^2}{\rho_0} + \frac{1}{K} \frac{(\theta_0 - 1)^2}{\theta_0} \right] dx$$
(2.18)

and taking into account (1.9), (1.10) and (2.5) we immediately get

$$U_n(0) \le C. \tag{2.19}$$

Now, multiply (1.1), (1.2), (1.3) and (1.4) by  $\rho^{-1}(1-\rho^{-1})$ ,  $K^{-1}v$ ,  $A^{-1}K^{-1}\omega\rho^{-1}$ and  $K^{-1}(1-\theta^{-1})\rho^{-1}$ , respectively, and integrate over ]-n, n[ and over  $]0, t[, t \in ]0, T[$ . After addition of the obtained equations we find that

$$U_n(t) + \int_0^t V_n(\tau) d\tau = U_n(0) \le C.$$
 (2.20)

**Lemma 2.1.** *For*  $t \in (0, T)$ ,

$$\|v_n(t)\|_{L^2(]-n,n[)} \le C , \qquad (2.21)$$

$$\|\omega_n(t)\|_{L^2(]-n,n[)} \le C.$$
(2.22)

**Proof.** These estimates follow from (2.20) and the inequalities

$$\frac{1}{2K} \|v_n(t)\|_{L^2(]-n,n[)}^2 \le U_n(t), \ \frac{1}{2AK} \|\omega_n(t)\|_{L^2(]-n,n[)}^2 \le U_n(t).$$
(2.23)

Like in [1, pp. 68-75] we can conclude that for each subset ]m, m + 1[,  $m \in \{-n, -n + 1, ..., n - 1\}$ , of ] - n, n[ there exists  $a_m(t) \in ]m, m + 1[$  such that the restriction of  $\rho_n$  to  $Q'_{mT} = ]m, m + 1[\times]0, T[$  has the form

$$\rho_n(x,t) = \frac{\rho_{0n}(x)Y_{nm}(t)B_{nm}(x,t)}{1 + K\rho_{0n}(x)\int_0^t Y_{nm}(\tau)B_{nm}(x,\tau)\theta_n(x,\tau)d\tau},$$
(2.24)

where

$$Y_{nm}(t) = \frac{1}{\rho_{0n}(a_m(t))} exp\left\{ K \int_0^t \rho_n(a_m(t), \tau) \theta_n(a_m(t), \tau) d\tau \right\} , \qquad (2.25)$$

$$B_{nm}(x,t) = \rho_n(a_m(t),t)exp\left\{\int_{a_m(t)}^x [v_{0n}(\xi) - v_n(\xi,t)]d\xi\right\}.$$
 (2.26)

Also, there exist constants  $C_i(i = 1, .., 5)$  (independent of m and n) such that the estimates

$$C_1 \le \int_m^{m+1} \theta_n(x,t) dx \le C_2 ,$$
 (2.27)

$$C_3^{-1} \le B_{nm}(x,t) \le C_3 \ , \ C_4 \le Y_{nm}(t) \le C_5$$
 (2.28)

are satisfied for  $t \in ]0, T[$  and  $(x, t) \in Q'_{mT}$ . Because of (2.14) and (2.15) there exist positive functions

$$m_{\rho_n}(t) = \inf_{x \in ]-n,n[} \rho_n(x,t) , \ m_{\theta_n}(t) = \inf_{x \in ]-n,n[} \theta_n(x,t),$$
(2.29)

$$M_{\rho_n}(t) = \sup_{x \in ]-n,n[} \rho_n(x,t) \ , \ M_{\theta_n}(t) = \sup_{x \in ]-n,n[} \theta_n(x,t)$$
(2.30)

defined on ]0, T[ and we have the following results.

**Lemma 2.2.** *For*  $t \in ]0, T[,$ 

$$M_{\rho_n}(t) \le C , \qquad (2.31)$$

$$m_{\rho_n}(t) \ge C \left(1 + \int_0^t M_{\theta_n}(\tau) d\tau\right)^{-1}.$$
(2.32)

**Proof.** Using (1.9), (2.11), (2.29), (2.30) and estimates (2.28) from (2.24) we get (2.31) and (2.32).  $\Box$ 

We define non-negative functions  $I_{1n}$  and  $I_{2n}$  in ]0, T[ by

$$I_{1n}(t) = \int_{-n}^{n} \rho_n(x,t) \left(\frac{\partial \theta_n}{\partial x}(x,t)\right)^2 dx , \qquad (2.33)$$

$$I_{2n}(t) = \int_0^t I_{1n}(\tau) d\tau .$$
 (2.34)

Obviously,  $I_{1n}$  and  $I_{2n}$  belong to  $L^1(]0, T[)$ .

**Lemma 2.3.** For  $\varepsilon > 0$  sufficiently small, there exists a constant  $C_{\varepsilon} \in \mathbb{R}^+$  such that, for  $t \in [0, T[$ , the inequality

$$M_{\theta_n}^2(t) \le \varepsilon I_{1n}(t) + C_{\varepsilon}(1 + I_{2n}(t))$$
(2.35)

 $holds\ true.$ 

**Proof.** We introduce the function  $\psi_{nm}$  on  $Q'_{mT}$  by

$$\psi_{nm}(x,t) = \theta_n(x,t) - \int_m^{m+1} \theta_n(x,t) dx .$$
 (2.36)

There exists  $x_m(t) \in ]m, m+1[$  such that  $\psi_{nm}(x_m(t), t) = 0$ . By means of the Hölder inequality we find that

$$|\psi_{nm}(x,t)|^{\frac{3}{2}} \leq \int_{x_m(t)}^x \frac{\partial}{\partial \xi} |\psi_{nm}(\xi,t)|^{\frac{3}{2}} d\xi \leq \frac{3}{2} (\int_m^{m+1} \rho_n^{-1}(\xi,t) |\psi_{nm}(\xi,t)| d\xi)^{\frac{1}{2}} (\int_m^{m+1} \rho_n(\xi,t) (\frac{\partial \psi_{nm}}{\partial \xi}(\xi,t))^2 d\xi)^{\frac{1}{2}}.$$
 (2.37)

Because of (2.27) we have  $\int_{m}^{m+1} |\psi_{nm}(\xi, t)| d\xi \leq C$  (independently of m and n). Taking into account (2.32), (2.27), (2.33) and  $\frac{\partial \psi_{nm}}{\partial \xi}(\xi, t) = \frac{\partial \theta_n}{\partial \xi}(\xi, t)$  from (2.37) we obtain

$$M_{\theta_n}^2(t) \le C\left[\left(1 + \int_0^t M_{\theta_n}(\tau)d\tau\right)^{\frac{2}{3}} (I_{1n}(t))^{\frac{2}{3}} + 1\right], \ t \in ]0, T[.$$
(2.38)

Applying the Young inequality with parameter  $\varepsilon > 0$ , from (2.38) it follows

$$M_{\theta_n}^2(t) \le \varepsilon I_{1n}(t) + C_{\varepsilon} \left( 1 + \int_0^t M_{\theta_n}^2(\tau) d\tau \right)$$
(2.39)

and by means of the Gronwall's inequality ([1, p.25]) we get (2.35).

Now, we introduce the function

$$\Phi_n = \frac{1}{2}v_n^2 + \frac{1}{2A}\omega_n^2 + (\theta_n - 1) \text{ on } Q_{nT} .$$
(2.40)

Multiply equations (1.2), (1.3) and (1.4) by  $v_n\Phi_n$ ,  $A^{-1}\rho_n^{-1}\omega_n\Phi_n$  and  $\rho_n^{-1}\Phi_n$ , respectively, and integrate over ] - n, n[ using (2.6) and (2.7). After addition of the obtained equations, we have

$$\frac{1}{2}\frac{d}{dt}\int_{-n}^{n}\Phi_{n}^{2}dx + \int_{-n}^{n}\rho_{n}\left(\frac{\partial\Phi_{n}}{\partial x}\right)^{2} + (1-A^{-1})\int_{-n}^{n}\rho_{n}\frac{\partial\omega_{n}}{\partial x}\omega_{n}\frac{\partial\Phi_{n}}{\partial x}dx + (D-1)\int_{-n}^{n}\rho_{n}\frac{\partial\theta_{n}}{\partial x}\frac{\partial\Phi_{n}}{\partial x}dx - K\int_{-n}^{n}\rho_{n}\theta_{n}v_{n}\frac{\partial\Phi_{n}}{\partial x}dx = 0$$
(2.41)

on ]0, T[. Taking into account (2.21), (2.31) and (2.35), in the same way as in [8, Lemma 2.4], we conclude that the inequality

$$\frac{d}{dt} \left( \int_{-n}^{n} (\Phi_n^2 + C_1 v_n^4 + C_2 \omega_n^4) dx + DI_{2n} \right) \\
\leq C \left( 1 + \int_{-n}^{n} (\Phi_n^2 + C_1 v_n^4 + C_2 \omega_n^4) dx + DI_{2n} \right)$$
(2.42)

holds true. Using the embedding  $H^1(R) \subset B^0(R)$ , (2.2), (1.10) and (2.5) we obtain that  $\|\Phi_n(0)\|_{L^2(]-n,n[)}^2 \leq C$  and after integration (2.42) over  $]0,t[, t \in ]0,T[$ , we find that

$$\int_{-n}^{n} (\Phi_n^2 + C_1 v_n^4 + C_2 \omega_n^4) dx + DI_{2n} \le C \text{ on } ]0T[.$$
(2.43)

**Lemma 2.4.** *For*  $t \in (0, T)$ ,

$$\|(\theta_n - 1)(t)\|_{L^2(]-n,n[)} \le C , \qquad (2.44)$$

$$\int_0^t M_{\theta_n}^2(\tau) d\tau \le C , \qquad (2.45)$$

$$m_{\rho_n}(t) \ge C , \qquad (2.46)$$

$$\int_{0}^{t} \left\| \frac{\partial \theta_{n}}{\partial x}(\tau) \right\|_{L^{2}(]-n,n[)}^{2} d\tau \leq C.$$
(2.47)

**Proof.** Estimate (2.44) follows from (2.43) and the inequality  $\|\theta_n - 1\|_{L^2(]-n,n[)}^2 \leq \int_{-n}^{n} \Phi_n^2 dx$ . Integrating (2.35) over ]0, t[ and taking into account (2.33), (2.34) and

(2.43) we get (2.45). From (2.32) and (2.45) we obtain (2.46). At last, using (2.46) and the estimate for  $I_{2n}$  from (2.43) we conclude that (2.47) holds.

Differentiating equality (2.24) with respect to x we get

$$\frac{\partial \rho_n}{\partial x} = \rho_n \varphi_n - \rho_n^2 Y_{nm}^{-1} B_{nm}^{-1} \left[ \frac{d}{dx} (\frac{1}{\rho_{0n}}) + K \int_0^t B_{nm} Y_{nm} (\frac{\partial \theta_n}{\partial x} + \theta_n \varphi_n) d\tau \right], (2.48)$$

where  $\varphi_n(x,t) = v_{0n}(x) - v_n(x,t)$ . Using (1.9), (1.10), (2.28) and (2.31) from (2.48) we obtain

$$\left\| \frac{\partial \rho_n}{\partial x}(t) \right\|_{L^2(]-n,n[)}^2 \leq C \left( \|v_{0n}\|_{L^2(]-n,n[)}^2 + \|v_n(t)\|_{L^2(]-n,n[)}^2 \right) + C \left[ 1 + \int_0^t \left\| \frac{\partial \theta_n}{\partial x}(\tau) \right\|_{L^2(]-n,n[)}^2 d\tau + \int_0^t M_{\theta_n}^2(\tau) \left( \|v_{0n}\|_{L^2(]-n,n[)}^2 + \|v_n(\tau)\|_{L^2(]-n,n[)}^2 \right) d\tau \right], \quad t \in ]0,T[.$$

$$(2.49)$$

**Lemma 2.5.** *For*  $t \in ]0, T[$ *,* 

$$\left\|\frac{\partial\rho_n}{\partial x}(t)\right\|_{L^2(]-n,n[)} \le C.$$
(2.50)

**Proof.** By means of estimates (2.21), (2.45), (2.47) and (2.5) the result follows directly from (2.49).

Multiplying (1.2) and (1.3) by  $v_n$  and  $\rho_n^{-1}\omega_n$ , respectively, integrating over ] - n, n[ and using (2.21) and (2.50) in the first equation, we find that

$$\frac{1}{2}\frac{d}{dt}\int_{-n}^{n}v_{n}^{2}dx + \int_{-n}^{n}\rho_{n}\left(\frac{\partial v_{n}}{\partial x}\right)^{2}dx \leq KM_{\theta_{n}}\left\|\frac{\partial\rho_{n}}{\partial x}\right\|_{L^{2}(]-n,n[)}\|v_{n}\|_{L^{2}(]-n,n[)} + C\left\|\frac{\partial\theta_{n}}{\partial x}\right\|_{L^{2}(]-n,n[)}\|v_{n}\|_{L^{2}(]-n,n[)} \leq C\left(M_{\theta_{n}} + \left\|\frac{\partial\theta_{n}}{\partial x}\right\|_{L^{2}(]-n,n[)}\right), \quad (2.51)$$

$$\frac{1}{2}\frac{d}{dt}\int_{-n}^{n}\omega_{n}^{2}dx + A\int_{-n}^{n}\rho_{n}\left(\frac{\partial\omega_{n}}{\partial x}\right)^{2}dx + A\int_{-n}^{n}\frac{\omega_{n}^{2}}{\rho_{n}}dx = 0 \text{ on } ]0,T[.$$
(2.52)

**Lemma 2.6.** For  $t \in [0, T[$ ,

$$\int_{0}^{t} \left\| \frac{\partial v_{n}}{\partial x}(\tau) \right\|_{L^{2}(]-n,n[)}^{2} d\tau \leq C,$$
(2.53)

$$\int_0^t \left\| \frac{\partial \omega_n}{\partial x}(\tau) \right\|_{L^2(]-n,n[)}^2 d\tau \le C.$$
(2.54)

**Proof.** Integrating (2.51) and (2.52) over  $]0, t[, t \in ]0, T[$ , and applying (2.45)-(2.47), (2.5) and (2.31) we get (2.53) and (2.54).

Now, we write equation (1.1) in the form

$$\frac{\partial}{\partial t} \left( \frac{1}{\rho_n} \right) = \frac{\partial v_n}{\partial x}.$$
(2.55)

Integrating over  $]0,t[,t\in]0,T[$  , squaring and integrating again over ]-n,n[, we obtain the inequality

$$\int_{-n}^{n} \left(\frac{1-\rho_n}{\rho_n}\right)^2 \le C \left[\int_{-n}^{n} \left(\frac{1-\rho_{0n}}{\rho_{0n}}\right)^2 dx + \int_{0}^{t} \int_{-n}^{n} \left(\frac{\partial v_n}{\partial x}\right)^2 dx d\tau\right].$$
 (2.56)

**Lemma 2.7.** *For*  $t \in ]0, T[$ *,* 

$$\|(\rho_n - 1)(t)\|_{L^2(]-n,n[)}^2 \le C.$$
(2.57)

**Proof.** Using (1.9), (1.10), (2.31) and (2.53), from (2.56) we easily get (2.57).  $\Box$ 

**Lemma 2.8.** For  $t \in [0, T[$ ,

$$\left\|\frac{\partial v_n}{\partial x}(t)\right\|_{L^2(]-n,n[)}^2 + \int_0^t \left\|\frac{\partial^2 v_n}{\partial x^2}(\tau)\right\|_{L^2(]-n,n[)}^2 d\tau \le C,$$
(2.58)

$$\left\|\frac{\partial\omega_n}{\partial x}(t)\right\|_{L^2(]-n,n[)}^2 + \int_0^t \left\|\frac{\partial^2\omega_n}{\partial x^2}(\tau)\right\|_{L^2(]-n,n[)}^2 d\tau \le C,\tag{2.59}$$

$$\left\|\frac{\partial\theta_n}{\partial x}(t)\right\|_{L^2(]-n,n[)}^2 + \int_0^t \left\|\frac{\partial^2\theta_n}{\partial x^2}(\tau)\right\|_{L^2(]-n,n[)}^2 d\tau \le C.$$
(2.60)

**Proof.** After multiplying (1.2) by  $\partial^2 v_n / \partial x^2$  and integrating by parts over ]-n, n[ and over ]0, t[, in the same way as in [1, pp.53-54], we obtain (2.58). Multiplying (1.3) and (1.4) by  $A^{-1}\rho_n^{-1}\partial^2\omega_n/\partial x^2$  and  $\rho_n^{-1}\partial^2\theta_n/\partial x^2$ , respectively, and integrating by parts over ]-n, n[ and over ]0, t[, in the same way as in [8, Lemmas 2.7, 2.8] we get estimates (2.59) and (2.60).

**Lemma 2.9.** *For*  $t \in (0, T)$ ,

$$\int_{0}^{t} \left\| \frac{\partial \rho_{n}}{\partial t}(\tau) \right\|_{L^{2}(]-n,n[)}^{2} d\tau \leq C,$$
(2.61)

$$\int_{0}^{t} \left\| \frac{\partial v_{n}}{\partial t}(\tau) \right\|_{L^{2}(]-n,n[)}^{2} d\tau \leq C, \qquad (2.62)$$

$$\int_{0}^{t} \left\| \frac{\partial \omega_{n}}{\partial t}(\tau) \right\|_{L^{2}(]-n,n[)}^{2} d\tau \leq C,$$
(2.63)

$$\int_0^t \left\| \frac{\partial \theta_n}{\partial t}(\tau) \right\|_{L^2(]-n,n[)}^2 d\tau \le C.$$
(2.64)

**Proof.** We square equations (1.1) and (1.2), integrate over ] - n, n[ and ]0, t[. Then in the same way as in [1, pp.53-54] we get (2.61) and (2.62). Also, squaring equations (1.3) and (1.4), integrating over ] - n, n[ and ]0, t[ in the same way as in [8, Lemmas 2.7, 2.8] we obtain (2.63) and (2.64).

# 3. Proof of Theorem 1.1

Let us denote again by  $\rho_n$  and  $\theta_n$  the extensions of  $\rho_n$  and  $\theta_n$  by 1 from  $Q_{nT}$  to  $\Pi$  and by  $v_n$  and  $\omega_n$  the extensions of  $v_n$  and  $\omega_n$  by zero outside of  $Q_{nT}$ .

We can find a function  $\varphi \in \mathcal{D}(R)$  such that

$$\varphi(x) = \begin{cases} 1 & \text{if } |x| \le 1\\ 0 & \text{if } |x| \ge 2 \end{cases}$$

$$(3.1)$$

and then we define  $\varphi_n$  by

$$\varphi_n(x) = \varphi(\frac{2x}{n}), \quad n \in N.$$
 (3.2)

For  $v_n$  and  $\omega_n$  we put

$$\bar{v}_n = v_n \varphi_n, \ \bar{\omega}_n = \omega_n \varphi_n$$
 (3.3)

and for  $\rho_n$  and  $\theta_n$  we introduce

$$\bar{\rho}_n = (\rho_n - 1)\varphi_n + 1, \quad \bar{\theta}_n = (\theta_n - 1)\varphi_n + 1. \tag{3.4}$$

One can easily conclude that the function  $\bar{S}_n = (\bar{\rho}_n, \bar{v}_n, \bar{\omega}_n, \bar{\theta}_n)$  satisfies system (1.1)-(1.4) a.e. in  $] - \frac{n}{2}, \frac{n}{2}[\times]0, T[$  and initial data (2.1), (2.2) and (2.4) a.e. in  $] - \frac{n}{2}, \frac{n}{2}[$ . Using the properties of  $\rho_n, v_n, \omega_n$  and  $\theta_n$  from (3.2)-(3.4) we observe that

$$\bar{\rho}_{n0} - 1 \rightarrow \rho_0 - 1, \ \bar{\theta}_{n0} - 1 \rightarrow \theta_0 - 1 \text{ strongly in } L^2(R)$$
 (3.5)

and

$$\bar{v}_{n0} \to v_0, \ \bar{\omega}_{n0} \to \omega_0 \text{ strongly in } L^2(R),$$
(3.6)

where

$$\bar{\rho}_{n0} = \bar{\rho}_n(x,0), \ \bar{\theta}_{n0} = \theta_n(x,0),$$
$$\bar{v}_{n0} = \bar{v}_n(x,0), \ \bar{\omega}_{n0} = \bar{\omega}_n(x,0), \ x \in \mathbb{R}.$$
(3.7)

In order to simplify a notation in what follows we write  $\rho_n$  instead  $\bar{\rho}_n$ , etc..

From Lemmas 2.5, 2.7 and 2.9 we conclude that

$$\{\rho_n - 1\}$$
 is bounded in  $L^{\infty}(0,T; H^1(R))$  and  $H^1(\Pi)$ . (3.8)

Moreover, taking into account (2.31) and (2.46), from (3.4) we obtain that

$$\{\rho_n\}$$
 is bounded in  $L^{\infty}(\Pi)$ . (3.9)

By means of Lemmas 2.1, 2.4, 2.6, 2.8 and 2.9 from (3.3) and (3.4) we get that

$$\{v_n\}, \{\omega_n\}, \{\theta_n - 1\}$$
 are bounded in  $L^{\infty}(0, T; H^1(R)), H^1(\Pi)$  (3.10)

and 
$$L^2(0,T; H^2(R))$$
.

Lemma 3.1. There exists a function

$$\rho - 1 \in H^1(\Pi) \cap L^{\infty}(0, T; H^1(R))$$
(3.11)

and a subsequence of  $\{\rho_n - 1\}$  (for simplicity denoted again as  $\{\rho_n - 1\}$ ) such that

$$\rho_n - 1 \to \rho - 1 \text{ weakly}^* \text{ in } L^{\infty}(0, T; H^1(R)),$$
(3.12)

$$\rho_n - 1 \to \rho - 1 \text{ weakly in } H^1(\Pi). \tag{3.13}$$

The function  $\rho$  belongs to  $L^{\infty}(\Pi)$  and has the properties:

$$\rho(x,0) = \rho_0(x) \ a.e. \ in \mathbb{R},$$
(3.14)

$$m_1 \le \rho \le M_1 \quad a.e. \ in \Pi, \tag{3.15}$$

where  $m_1, M_1 \in \mathbb{R}^+$ .

**Proof.** Since the sequence  $\{\rho_n - 1\}$  is bounded in  $L^{\infty}(0, T; H^1(R))$  (dual of  $L^1(0, T; H^{-1}(R))$ ), it is possible to extract a subsequence (denoted again as  $\{\rho_n - 1\}$ ) such that  $\rho_n - 1 \rightarrow \rho - 1$  weakly<sup>\*</sup> in  $L^{\infty}(0, T; H^1(R))$  (see [4, pp.498-503]). It means that for  $g \in L^1(0, T; H^{-1}(R))$ ,  $(g(t) = (g_1(t), g_2(t)) \in L^2(R) \times L^2(R))$  we have

$$\int_{\Pi} (\rho_n - 1)g_1 dx dt + \int_{\Pi} \frac{\partial \rho_n}{\partial x} g_2 dx dt \to \int_{\Pi} (\rho - 1)g_1 dx dt + \int_{\Pi} \frac{\partial \rho}{\partial x} g_2 dx dt .$$
(3.16)

Specially, for all  $\varphi \in \mathcal{D}(\Pi)$  from (3.16) we obtain

$$\int_{\Pi} (\rho_n - 1)\varphi dx dt \to \int_{\Pi} (\rho - 1)\varphi dx dt , \qquad (3.17)$$

$$\int_{\Pi} \frac{\partial \rho_n}{\partial x} \varphi dx dt \to \int_{\Pi} \frac{\partial \rho}{\partial x} \varphi dx dt .$$
(3.18)

Also,  $\{\rho_n\}$  is bounded in  $L^{\infty}(\Pi)$  and therefore there exists a subsequence (denoted by  $\{\rho_n\}$ ) such that  $\rho_n \to \rho$  weakly<sup>\*</sup> in  $L^{\infty}(\Pi)$ . Specially, for all  $\varphi \in \mathcal{D}(\Pi)$  we get

$$\int_{\Pi} \rho_n(x,t)\varphi(x,t)dxdt \to \int_{\Pi} \rho(x,t)\varphi(x,t)dxdt .$$
(3.19)

Because of (3.8) we can take a further subsequence of  $\{\rho_n - 1\}$  such that  $\rho_n - 1 \rightarrow \rho - 1$  weakly in  $H^1(\Pi)$ . From this convergence we find out that for  $\varphi \in \mathcal{D}(\Pi)$ , it holds

$$\int_{\Pi} \frac{\partial \rho_n}{\partial t}(x,t)\varphi(x,t)dxdt \to \int_{\Pi} \frac{\partial \rho}{\partial t}(x,t)\varphi(x,t)dxdt .$$
(3.20)

Statement (3.11) is a consequence of the above convergences.

Taking into account (2.31), (2.46), (3.1), (3.2), (3.4) and (3.19) we conclude that there exist  $m_1, M_1 \in \mathbb{R}^+$  such that (3.15) holds. From the embedding theorem (see [4, p.473]) we observe that functions  $\rho_n - 1, \rho - 1$  belong to  $C([0, T]; L^2(R))$  being equipped with the norm of uniform convergence. Now we may speak of the traces  $\rho_n(x, 0) - 1$  and  $\rho(x, 0) - 1$ .

Let  $\psi \in C^{\infty}([0,T]), \psi(0) \neq 0$  and  $\psi$  vanishes in a neighbourhood of T. Applying Green's formula ([4, p.477]) we obtain

$$\int_0^T \int_R \frac{\partial \rho_n}{\partial t}(x,t)u(x)\psi(t)dxdt + \int_0^T \int_R (\rho_n - 1)(x,t)u(x)\frac{d\psi}{dt}(t)dxdt$$
$$= -\psi(0)\int_R (\rho_{n0} - 1)u(x)dx , \quad (3.21)$$

$$\int_0^T \int_R \frac{\partial \rho}{\partial t}(x,t)u(x)\psi(t)dxdt + \int_0^T \int_R (\rho-1)(x,t)u(x)\frac{d\psi}{dt}(t)dxdt$$
$$= -\psi(0)\int_R (\rho(x,0)-1)u(x)dx , \qquad (3.22)$$

for all  $u \in \mathcal{D}(R)$ . Comparing (3.21) and (3.22) (when  $n \to \infty$ ) and using (3.17), (3.20) and (3.5) we find that  $\rho(x, 0) = \rho_0(x)$  in the sense of distributions in  $\mathbb{R}$ .  $\Box$ Lemma 3.2. There exist functions

$$v, \omega, \theta - 1 \in L^{\infty}(0, T; H^1(R)) \cap H^1(\Pi) \cap L^2(0, T; H^2(R))$$
 (3.23)

and a subsequence of  $\{v_n, \omega_n, \theta_n - 1\}$  (denoted again as  $\{v_n, \omega_n, \theta_n - 1\}$ ) such that

$$\{v_n, \omega_n, \theta_n - 1\} \to \{v, \omega, \theta - 1\}$$
 weakly \* in  $(L^{\infty}(0, T; H^1(R))^3,$  (3.24)

$$\{v_n, \omega_n, \theta_n - 1\} \to \{v, \omega, \theta - 1\} \quad weakly \ in \ (H^1(R))^3 , \qquad (3.25)$$

$$\{v_n, \omega_n, \theta_n - 1\} \to \{v, \omega, \theta - 1\}$$
 weakly in  $(L^2(0, T; H^2(R))^3)$ . (3.26)

Functions  $v, \omega$  and  $\theta$  have the properties:

$$v(x,0) = v_0(x) , \ \omega(x,0) = \omega_0(x) , \ \theta(x,0) = \theta_0(x) \ a.e. \ in \mathbf{R}.$$
 (3.27)

**Proof.** Conclusions (3.23)-(3.26) follow immediately from (3.10). From the weak convergences we conclude that for  $\varphi \in \mathcal{D}(\Pi)$ , it follows

$$\int_{\Pi} v_n(x,t)\varphi(x,t)dxdt \to \int_{\Pi} v(x,t)\varphi(x,t)dxdt, \qquad (3.28)$$

$$\int_{\Pi} \frac{\partial v_n}{\partial x}(x,t)\varphi(x,t)dxdt \to \int_{\Pi} \frac{\partial v}{\partial x}(x,t)\varphi(x,t)dxdt,$$
(3.29)

$$\int_{\Pi} \frac{\partial v_n}{\partial t}(x,t)\varphi(x,t)dxdt \to \int_{\Pi} \frac{\partial v}{\partial t}(x,t)\varphi(x,t)dxdt,$$
(3.30)

$$\int_{\Pi} \frac{\partial^2 v_n}{\partial^2 x}(x,t)\varphi(x,t)dxdt \to \int_{\Pi} \frac{\partial^2 v}{\partial^2 x}(x,t)\varphi(x,t)dxdt$$
(3.31)

(when  $n \to \infty$ ), which is true for  $\{\omega_n\}$  and  $\{\theta_n - 1\}$  also. By means of Green's formula we get properties (3.27) in the same way as (3.14).

**Lemma 3.3.** Functions  $\rho, v, \omega$  and  $\theta$ , defined by Lemma 3.1 and Lemma 3.2 satisfy equations (1.1)-(1.4) a.e. in  $\Pi$ .

**Proof.** Let  $\{S_n = (\rho_n, v_n, \omega_n, \theta_n) : n \in \mathbf{N}\}$  be the subsequence defined by *Lemmas 3.1* and *3.2*. By means of (3.9) and (3.15) we obtain the inequalities

$$\left| \int_{\Pi} (\rho_n^2 \frac{\partial v_n}{\partial x} - \rho^2 \frac{\partial v}{\partial x}) \varphi dx dt \right| \leq \left| \int_{\Pi} \rho_n^2 (\frac{\partial v_n}{\partial x} - \frac{\partial v}{\partial x}) \varphi dx dt \right| \\ + \left| \int_{\Pi} \frac{\partial v}{\partial x} (\rho_n - \rho) (\rho_n + \rho) \varphi dx dt \right| \\ \leq C \left| \int_{\Pi} (\frac{\partial v_n}{\partial x} - \frac{\partial v}{\partial x}) \varphi dx dt \right| \\ + C \left| \int_{\Pi} \frac{\partial v}{\partial x} (\rho_n - \rho) \varphi dx dt \right| , \qquad (3.32)$$

for all  $\varphi \in \mathcal{D}(\Pi)$  and after integrating by parts we get

$$\left| \int_{\Pi} (\rho_n^2 \frac{\partial v_n}{\partial x} - \rho^2 \frac{\partial v}{\partial x}) \varphi dx dt \right| \leq C \left| \int_{\Pi} (\frac{\partial v_n}{\partial x} - \frac{\partial v}{\partial x}) \varphi dx dt \right| + C \left| \int_{\Pi} v (\frac{\partial \rho_n}{\partial x} - \frac{\partial \rho}{\partial x}) \varphi dx dt \right| + C \left| \int_{\Pi} v (\rho_n - \rho) \frac{\partial \varphi}{\partial x} dx dt \right| .$$
(3.33)

Taking into account (3.29), (3.19), (3.18) and (1.18), we conclude that for all  $\varphi \in \mathcal{D}(\Pi)$ , from (3.33) it follows

$$\int_{\Pi} \rho_n^2 \frac{\partial v_n}{\partial x} \varphi dx dt \to \int_{\Pi} \rho^2 \frac{\partial v}{\partial x} \varphi dx dt.$$
(3.34)

For  $\varphi \in \mathcal{D}(\Pi)$  there exists  $n_0 \in \mathbb{N}$  such that for all  $n \geq n_0$  functions  $\rho_n$  and  $v_n$  satisfy (1.1) in the sense of distributions in  $\Pi$  and therefore we have

$$\int_{\Pi} \left( \frac{\partial \rho_n}{\partial t} + \rho_n^2 \frac{\partial v_n}{\partial x} \right) \varphi dx dt = 0.$$
(3.35)

Applying (3.20) and (3.34), from (3.35) we get

$$\frac{\partial \rho}{\partial t} + \rho^2 \frac{\partial v}{\partial x} = 0$$
 a.e in II. (3.36)

In the same way one can prove that equations (1.2), (1.3) and (1.4) are satisfied.  $\Box$ 

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