

## Electronic Structure of Superheavy Elements and Their Compounds

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To determine if element 104 in its compounds has d- or p-character the electronic structure of  $KuCl_4$  has been calculated using a self-consistent relativistic multiple scattering model with  $X\alpha$ -exchange. By comparing the data obtained for  $KuCl_4$  with those for  $HfCl_4$  and  $PbCl_4$  it has been established that  $KuCl_4$  is not a full analog of  $HfCl_4$ .

### 1. INTRODUCTION

Long before the discovery of trans-actinides, questions were asked about the characted of the elements begining with N 103. After the discovery of Lr, the first experiment showed<sup>1</sup> that in water solution this element formed a stable trivalent state. By analogy with Lu its ground state configuration was supposed to be  $(Rn)f^{14}ds^2$ . In accordance with the fact the elements 103 to 112 were believed to form 6d-series, elements 113 to 118 were taken to form 7p-series. It meant that Lr had to be an analog of the elements of Ga-subgroup, and E104-Ti subgroup.

The following elements, 119 to 120, were considered to be analogs of alkaline and alkaline earth elements with ground state electronic configurations  $No/6d^{10}7p^68s^1$  and  $No/6d^{10}7p^68s^2$ . By tradition (La and Ac) element 121 was considered to have the electronic configuration  $No/6d^{10}7p^68s^27d$ , and then 32 electrons are to fill  $6f^{14}$  and  $5g^{18}$ -shells up to element 153. Further, there must be a 7d- and 8p-electron shell filling for elements 154 to 162 and 163 to 168. This filled-shell concept doesn't assume any deviations from the periodicity and from what would be expected from a straightforward extrapolation of the Periodic Table.

A number of atomic calculations that appeared in recent years have enabled an insight into the physical and chemical properties of the still unknown superheavy elements. Some results have been reviewed by Hermann<sup>3</sup> and Fricke<sup>4</sup>. Many authors (3—5) have also emphasized that further development in this field must be based on calculations (in this case relativistic adequately describing the electronic structure of molecules, containing superheavy elements. A better knowledge of their chemical and physical properties is the prerequisite for a successful extraction and identification of the possible artificially generated superheavy elements.

It is known that close to the end of the Periodic Table relativistic effects begin to play a more important role influencing the electronic structure of heavy atoms and molecules, and as a result the chemistry of their compounds. Nevertheless, up to now no element is known whose ground state is different compared to the normal continuation of the Periodic Table, due to the direct relativistic influence. The first candidate for such a deviation turned out to be element 103 (Lr). It was Mann<sup>6</sup> who first pointed out that single configuration Dirac-Fock calculations give  $5f^{14}7s^27p^1$  as the ground state. Independently Brewer<sup>7</sup> came to the same conclusion using a semi-empirical extrapolation and later he reaffirmed the result taking into account relativistic and correlation effects. Multi-configurational relativistic Dirac-Fock calculations<sup>9</sup> predict the same ( $5f^{14}7s^27p_{1/2}$ ) electronic ground state configuration which arises from a strong relativistic stabilization of the  $7p_{1/2}$  orbital. Furthermore, it has been suggested that the  $7s^2$  closed shell might be sufficiently stabilized by relativistic effects to make it an »inert core« so that only  $7p_{1/2}$  (or  $6d$ ) electron may be removed under reducing conditions, thus producing monovalent Lr. However, the experiments carried out in Darmstadt<sup>11</sup> have not confirmed the existence of monovalent Lr (it turned out to be impossible to reduce  $Lr^{3+}$  to the monovalent state in aqueous solution).

Experimental attempts to search for Lr as a volatile p-element<sup>12</sup> using on-line gas chromatography have not given any results so far. These experiments, however, do not allow the authors to exclude the existence of Lr(p). They believe that maybe in the presence of the chromatography column surface, due to the higher adsorption enthalpy of Lr(d) (750 kJ/mol) with respect to Lr(p) (—187 kJ/mol), the  $/Rn/5f^{14}7s^27p_{1/2}$  configuration promotes to  $/Rn/5f^{14}6d7s^2$ , which is predicted to be only about 93 kJ/mol higher in energy.

Taking into account the difficulties encountered by experimentators, the theoretical predictions in this area remain urgent.

### *Analogy on the Basis of Atomic Properties*

Considering Lr as a p-element one can hardly relate it to the elements of the Ga-subgroup, whose peculiarity is a high promotion energy to d-state. Thus, Tl which is very often compared with Lr has the first excited state  $6s^27s^2$  at  $26478\text{ cm}^{-1}$  above the ground state  $6s^26p_{1/2}$  ( $^2P_{1/2}$ ) and excitation energy to the d-state is even higher —  $36118\text{ cm}^{-1}$ . Lr differs from the typical p-elements both in character of excitation and in its energy. The ground state configuration of Lr ( $7f^{14}7s^27p$ ), as it has been shown<sup>9</sup>, has an admixture (24%) of  $7s^26d$  and  $6d7s7p$ -electronic states.

But Lr also differs from typical d-elements, *e.g.* the lowest excited levels (states) of Sc arise from the excitation of  $4s$ -electron to  $3d$ -level ( $3d^24s$ ) with energies of the transition being equal to  $11500$  and  $15000\text{ cm}^{-1}$  for  $^4F$  and  $^2F$ -states, respectively.

Earlier in the work<sup>13</sup> it was emphasized that promotion energies  $5f^97s^2 \rightarrow 5f^97s7p$  for the elements at the end of the actinide series are lower than those for  $5f^97s^2 \rightarrow 5f^97s6d$  excitation (Table I).

TABLE I  
Excitation energy  $f^i s^2 \rightarrow f^i sp$  and  $f^i s^2 \rightarrow f^i ds$ , in kcal/mol

	Bk	Cf	Es	Fm	Md	No
$s^2 \rightarrow sp$	49	51	54	55	57	59
$s^2 \rightarrow sd$	54	60	65	74	80	90

It is seen (Table I) that there is a following tendency to the end of actinide series: stabilization of 7p and destabilization of 6d-atomic orbitals (AO). This qualitative effect transforms into a quantitative one that takes place in the case of Lr.

Summing up all the experimental and theoretical data obtained, one can conclude that Lr might be an element with a new quality, and all experimental attempts to obtain it in a pure d or p-state hardly seem to be successful.

As far as E104 is concerned, it is supposed to begin a new series of elements in the Periodic Table, where 6d-shell is to fill with electrons. Data obtained by Zvara<sup>16</sup> about the properties of gaseous Ku halogenides and those obtained by Silva on the basis of the ion exchange behaviour of Ku are used as an evidence of the d-character of Ku. But the increasing role of relativistic effects must probably influence the electronic structure of Ku and its compounds. Keller<sup>14</sup> supposed that under the influence of relativistic stabilization of  $7p_{1/2}$ -AO the ground state configuration might be  $7s^2 7p_{1/2}^2$  and by its properties Ku might be an analog of Pb( $6s^2 6p^2$ ) and metallic Ku might be a volatile substance. But the lowest limit for the sublimation enthalpy for Ku  $\Delta H_s = 370$  kJ/mol and for Pb  $\Delta H_s = 190$  kJ/mol and in contrast to atomic Pb, Ku does not pass through the chromatography column in the hydrogen current<sup>15</sup>. But no final decision about the p-character of Ku could be reached in terms of the used experimental technique because it was impossible to determine the difference between  $\Delta H_s$  (Ku) and  $\Delta H_s$  (Hf).

Again, if one considers Ku as an analog of Pb, the latter has the first excited  $6s^2 6p 7s$ -state at  $35500 \text{ cm}^{-1}$  above the ground state, and the lowest d-state ( $6s^2 6p 6d$ ) at  $38800 \text{ cm}^{-1}$ . Nevertheless, the first excited state  $6s 5d^3$  of d-elements (Hf) is  $14500 \text{ cm}^{-1}$  above the ground state  $6s^2 5d^2$ . Relativistic Hartree-Fock-Slater and Dirac-Fock calculations<sup>4,13,18</sup> have given a number of close-spaced electronic configurations ( $7s^2 6d^2$ ,  $7s^2 6d 7p$ ,  $7s^2 7p^2$ ). Recent multiconfigurational Dirac-Fock calculations<sup>19</sup> predict that the ground state of Ku is the  $J = 2$  state with the main contribution from  $7s^2 6d 7p^2$  single configuration:

$$80\% 7s^2 6d 7p + 18\% 7s 7p 6d^2 + 2\% \text{ others.}$$

The first excited state  $6d^2 7s^2$  (95%) at about 0.5 eV is the  $J = 2$  state. The ground state of  $\text{Ku}^+$  is  $7s^2 6d$  by analogy to Hf, and the  $7s^2 7p$  configuration is 2.4 eV higher.

However, the question arises whether the fact that the ground state of Ku is  $7s^2 6d 7p$  instead of  $7s 6d^2$  is a good reason for believing that its properties would move to the p-elements. In accordance with the Huckel approach the strength of the chemical bonding formed by the atom in

compound is determined by the energies of valent orbitals and their space expansion and therefore the authors<sup>19</sup> consider Ku to be a typical d-element (AO energies of Ku are similar to Hf as well as values  $\langle r \rangle$ ).

The properties (p-or d-) that Ku will have in its compounds can be determined only on the basis of molecular calculations. In a sense,  $\text{KuCl}_4$  is the most proper species for both d- and p-elements form rather stable tetrachlorides.

### Method of Calculations and Parameters

The self-consistent (SCF) calculations were performed using a version of the multiple scattering  $X\alpha$ -method, a detailed description of which can be found elsewhere, *e. g.* in ref. 20.

In the radial Schrödinger-like equation the »large« radial function  $G_{nl}$  has the solution:

$$[H_{\text{NR}} + H_m + BH_D] G_{nl} = EG_{nl}$$

Here  $H_{\text{NR}}$  is the usual nonrelativistic operator for the central field problem,  $H_m$  and  $H_D$  are the standard massvelocity and Darwin operators and  $B$  is given by

$$B = \left\{ 1 + \frac{1}{4} \alpha^2 [E - V(r)] \right\}^{-1}$$

where  $\alpha = 1/137$ . The spin-orbital interaction is taken into account by using the perturbation theory:  $H'_{\text{SO}} = BH_{\text{SO}}$ , where  $H_{\text{SO}}$  is the standard spin-orbit operator of the Pauli theory. The collection of one-electron energies, obtained by adding the first-order perturbation theory shifts induced by the spin-orbit operator, will be labeled  $\text{RX}_\alpha\text{-SO}$  while the (more accurate) collection obtained by diagonalizing this operator among a group of adjacent levels will be labeled  $\text{RX}_\alpha\text{-SO}'$ . The latter procedure takes into account the effect of mixing several close-spaced electronic states for valent orbitals, which becomes more prominent for the heavy elements.

The bond length  $\text{Ku-Cl}$  was chosen as a sum of ionic radii (4.9905 a. u.) and we employed a  $\text{Hf-Cl}$  internuclear distance in  $\text{HfCl}_4$  of 4.8393 a. u. and 4.4045 a. u., and a  $\text{Pb-Cl}$  distance in  $\text{PbCl}_4$  of 4.5936 a. u.

We used the model of overlapping atomic spheres with their ratio being determined by the Norman procedure. The value of overlapping changed from 10 to 25%, the virial ratio being the best for 25% (in case  $\text{KuCl}_4$ ). SCF molecular calculations were carried out starting with atomic ground state configurations  $7s6d^2$ ,  $7s^27p^2$  and  $7s^27p6d$ .

### Results of Molecular Calculations and Discussion

It is well-known that the molecular field splitting of d- and p-orbitals is different. In the  $T_d$  field d-orbitals are split into  $e_g$ - and  $t_{2g}$  components and such a scheme is valid for all d-element compounds. But observing the molecular orbital (MO) scheme of any p-element compounds one can notice a peculiarity: a position of the first antibonding level of  $a_1$ -symmetry below  $t_2$ -level of pure p-character (Figure 1).

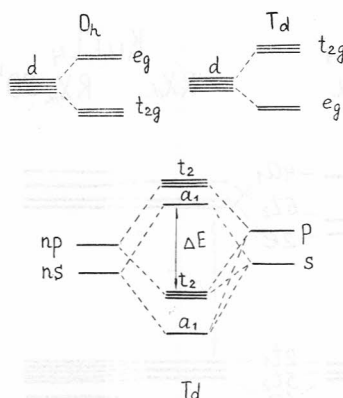


Figure 1. Principal scheme of d- and p-orbital splitting in the  $T_d$  molecular field.

Table II contains eigenvalues (MO energies) for  $HfCl_4$ ,  $KuCl_4$  and  $PbCl_4$ . Figure 2 combines this tabular information with the arrow indicating the intervening energy gap. It is obvious that the energy level orderings for  $HfCl_4$  and  $KuCl_4$  are similar except for the almost degenerate lower  $2a_1-1t_1$  MOs. Eigenvalues for  $HfCl_4$  and  $KuCl_4$  practically do not differ either. A more prominent distinction is in the antibonding MO ordering. Thus, in  $HfCl_4$  the classical scheme of d-orbitals splitting in  $T_d$ -crystal field is preserved, whereas in  $KuCl_4$   $4a_1$ -antibonding level dramatically decreases being the lowest unoccupied (LU) MO in the case of 25% overlapping (Figure 3). Figure 2 and Table II show that such a situation takes place in the MO scheme for  $PbCl_4$  where  $3a_1$  level is the LUMO. Thus, the electron structure of  $KuCl_4$  has an element peculiar for the p-element compounds.

TABLE II  
Eigenvalues ( $RXa$ ), eV

MO	$HfCl_4$ 25%	$KuCl_4$		MO	$PbCl_4$ 25%
		20%	25%		
$1t_2$	39.05	35.87	35.51		
$1a_1$	20.98	20.90	20.99		
$2t_2$	20.86	20.72	20.80	$1t_2$	21.86
$2a_1$	18.91	19.19	18.97	$1e$	20.61
$3t_2$	18.90	19.21	18.98	$1a_1$	21.47
$1t_1$	18.90	19.22	18.98	$2t_2$	19.82
$3a_1$	10.64	10.62	10.82	$2a_1$	13.58
$4t_2$	9.90	9.61	9.83	$3t_2$	9.80
$1e$	8.99	8.86	8.84	$2e$	8.58
$5t_2$	8.78	8.70	8.63	$4t_2$	8.54
$2t_1$	8.52	8.50	8.48	$1t_1$	8.46
$2e$	3.92	3.06	2.70	$3a_1$	4.82
$6t_2$	3.07	2.30	2.00	$5t_2$	0.93
$4a_1$	1.97	2.94	2.80	$4a_1$	0.32

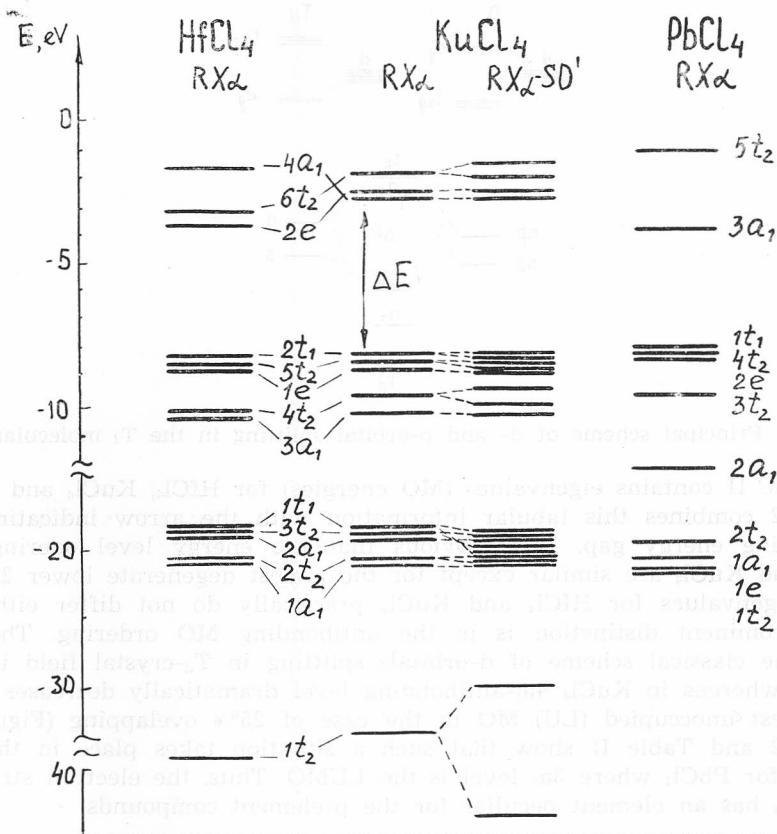


Figure 2. Energy level scheme (eigenvalues) for  $\text{HfCl}_4$ ,  $\text{KuCl}_4$  and  $\text{PbCl}_4$ .

The upper valence level ordering for  $\text{PbCl}_4$  is like that for  $\text{KuCl}_4$  ( $2a_1$ — $1t_1$ ) with similar eigenvalues. The lower valence levels for  $\text{PbCl}_4$  ( $1t_2$ — $2t_2$ ) differ from  $\text{KuCl}_4$  by either their ordering or eigenvalues. Besides, there is a MO of  $1e$ -symmetry (100% d) in  $\text{PbCl}_4$ . Looking for some analogy one can believe that  $4f^{14}$  and  $5f^{14}$  — inner shells in  $\text{HfCl}_4$  and  $\text{KuCl}_4$  (lower valence levels  $2a_1$ — $1t_1$ ) influence to the same extent the electronic structure of upper valence levels whereas the  $5d^{10}$ -shell in  $\text{PbCl}_4$  influences the upper valence levels to a different extent ( $1t_2$  and  $1e$  levels of  $5d$  (100%) character lie by 2 and 3 eV lower  $2a_1$ — $1t_1$ -levels of  $f$ (100%) character in  $\text{KuCl}_4$  and  $\text{HfCl}_4$ ).

Before analyzing the MO composition let us consider the influence of relativistic effects on the electronic structure of the species under consideration. For  $\text{HfCl}_4$  such effects as mass-velocity dependence and Darwin interaction do not change the energy level ordering in going from  $X_\alpha$  to  $\text{RX}_\alpha$  calculation. There is only a slight upward shift of unoccupied levels of d-character owing to an indirect relativistic effect leading to the destabilization of d-orbitals. As a result, the energy gap increases from 3.67 to 4.6 eV.

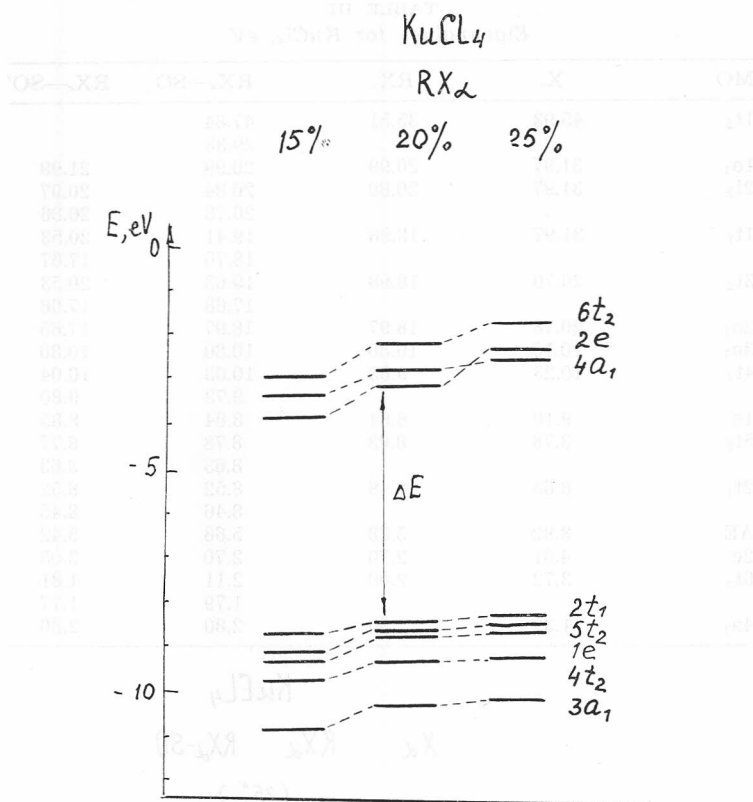


Figure 3. Upper bonding and low antibonding molecular orbitals for  $\text{KuCl}_4$  as a function of overlapping of atomic spheres.

In the case of  $\text{PbCl}_4$  there is also a slight shift of unoccupied levels in going from  $X_\alpha$  to  $\text{RX}_\alpha$  calculation, resulting, however, in a lowering of the LU  $3a_1$ -MO of s-character as a consequence of relativistic stabilization of 6s-orbitals. In this case the energy gap decreases from 4.58 to 3.64 eV. Spin-orbit interaction only slightly splits  $t_1$  and  $t_2$ -MO.

Table III (Figure 4) contains the results of all stages of calculations for  $\text{KuCl}_4$ . Relativistic effects ( $\text{RX}_\alpha$ -calculation) result in destabilization of d-orbitals ( $2e$  and  $6t_2$ ) and stabilization of  $4a_1$ -MO of 7s-character so that  $4a_1$ -MO becomes LUMO. Spin-orbit interaction ( $\text{RX}_\alpha$ -SO) does not influence the energy level ordering. The more accurate  $\text{RX}_\alpha$ -SO' procedure results in the LUMO being  $t_2$ -MO of mixed d-p-character. Thus, both  $\text{RX}_\alpha$ -SO and  $\text{RX}_\alpha$ -SO' calculations show that in  $\text{KuCl}_4$  the classical scheme of d-orbital splitting in the  $T_d$ -molecular field is not observed.

Table IV gives the charge composition (%) of the MOs. For bonding levels ( $3a_1$ — $2t_1$ ) d-orbitals of Ku participate to a greater extent in the chemical bond formation than p-orbitals ( $4t_2$ -MO), whereas in  $\text{PbCl}_4$  this orbital ( $3t_2$ ) is characterized by strong participation of p-AO and in the region of bonding orbitals ( $2a_1$ — $1t_1$ ) the s- and p-orbitals are predominant.

TABLE III  
Eigenvalues for  $KuCl_4$ , eV

MO	$X_\alpha$	$RX_\alpha$	$RX_\alpha-SO$	$RX_\alpha-SO'$
$1t_2$	45.93	35.51	47.84	
			29.33	
$1a_1$	31.97	20.99	20.99	21.99
$2t_2$	31.97	20.80	20.84	20.97
			20.78	20.86
$1t_1$	31.97	18.98	19.41	20.53
			18.76	17.67
$3t_2$	20.70	18.98	19.63	20.53
			17.68	17.66
$2a_1$	20.78	18.97	18.97	17.65
$3a_1$	10.10	10.80	10.80	10.80
$4t_2$	10.23	9.83	10.03	10.04
			9.73	9.80
$1e$	9.10	8.84	8.84	8.85
$5t_2$	8.78	8.63	8.78	8.77
			8.63	8.63
$2t_1$	8.63	8.48	8.52	8.52
			8.46	8.45
$\Delta E$	3.82	5.68	5.66	5.42
$2e$	4.81	2.70	2.70	3.03
$6t_2$	3.72	2.00	2.11	1.81
			1.79	1.77
$4a_1$	1.32	2.80	2.80	2.80

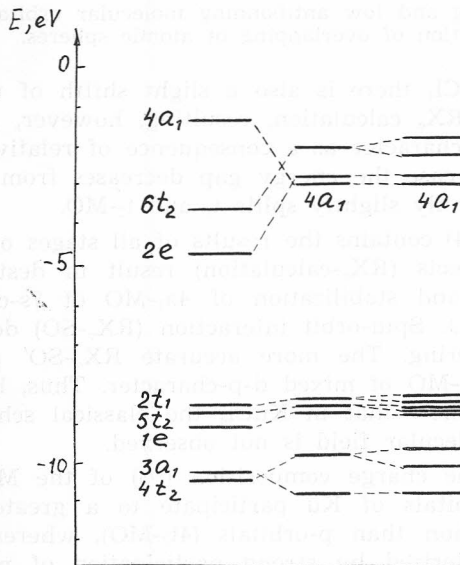
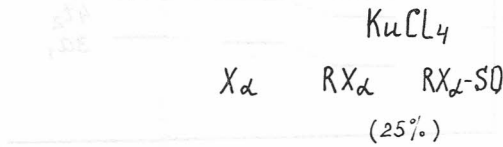


Figure 4. MO scheme for  $KuCl_4$  displaying the strong influence of relativistic effects on the unoccupied ( $2e-4a_1$ ) levels.



TABLE IV

MO composition (%). Only the predominant orbital contributions are given

MO	AO	HfCl <sub>4</sub>	KuCl <sub>4</sub>	MO	AO	PbCl <sub>4</sub>
1a <sub>1</sub>	f	10.59	12.96	1e	d	99.98
	s (Cl)	85.99	83.48	1a <sub>1</sub>	s	7.96
2t <sub>2</sub>	f	3.30	4.32		f	1.10
	s (Cl)	93.21	92.50		s (Cl)	89.60
2a <sub>1</sub>	f	99.70	99.94	2t <sub>2</sub>	p	2.56
3t <sub>2</sub>	f	99.90	99.98		d	2.64
1t <sub>1</sub>	f	99.99	100.00		s (Cl)	94.36
3a <sub>1</sub>	s	21.53	26.64	2a <sub>1</sub>	s	43.54
	f	1.70	2.02		f	0.58
	p (Cl)	74.80	69.07		p (Cl)	48.23
4t <sub>2</sub>	p	4.76	5.26	3t <sub>2</sub>	p	17.10
	d	16.41	13.51		d	1.72
	p (Cl)	77.94	79.80		p (Cl)	79.85
1e	d	9.69	7.24	2e	d	2.02
	p (Cl)	90.31	92.76		p (Cl)	97.98
5t <sub>2</sub>	p	2.86	2.61	4t <sub>2</sub>	p	0.02
	d	1.24	0.94		d	2.59
	p (Cl)	95.78	96.30		p (Cl)	96.53
2t <sub>1</sub>	p (Cl)	98.80	98.55	1t <sub>1</sub>	p (Cl)	98.57
2e	d	90.89	92.39			
	p (Cl)	9.11	7.69			
6t <sub>2</sub>	p	7.93	10.28	5t <sub>2</sub>	p	39.21
	d	75.13	74.71		d	8.19
	p (Cl)	12.61	10.85		p (Cl)	12.53
4a <sub>1</sub>	s	35.22	43.02	3a <sub>1</sub>	s	32.59
	s (Cl)	40.43	24.42		s (Cl)	70.64
	p (Cl)	14.54	22.99		p (Cl)	47.39

Comparing HfCl<sub>4</sub> and KuCl<sub>4</sub> one can notice that in the latter molecule the chemical bond is weaker due to the relativistic destabilization of the 6d-orbital, but at the same time it becomes stronger by the relativistic stabilization of p-orbitals (e. g. 4t<sub>2</sub>-MO). The highest occupied MO in all species is similar in composition (~99% 3s(Cl)).

For antibonding levels the MO composition for KuCl<sub>4</sub> is like that for HfCl<sub>4</sub> but with a slightly prominent p-character (6t<sub>2</sub>-MO).

Table V summarizes the results of AO population analysis. One can see that both in HfCl<sub>4</sub> and KuCl<sub>4</sub> 6p- and 7p-AOs have a similar electronic

TABLE V

AO populations for HfCl<sub>4</sub>, KuCl<sub>4</sub> and PbCl<sub>4</sub> (25% overlapping)

AO	HfCl <sub>4</sub>	AO	KuCl <sub>4</sub>	AO	PbCl <sub>4</sub>
6s	0.49	7s	0.59	6s	1.03
5p + 6p	6.49	6p + 7p	6.50	6p	1.18
5d	1.55	6d	1.26	5d + 6d	10.33
4f + 5f	14.54	5f + 6f	14.70	5f	0.21
Charges on metal atom and ligand (Cl)					
q <sub>Me</sub>	0.93	q <sub>Me</sub>	0.95	q <sub>Me</sub>	1.24
q <sub>L</sub>	-0.23	q <sub>L</sub>	-0.24	q <sub>L</sub>	-0.31

charge. This is quite reasonable because from atomic calculations it follows that the 5d-6p (2.95 eV) and 6d-7p (2.57 eV) energy differences in Hf and Ku, respectively, are similar. In PbCl<sub>4</sub> the 6d-orbital spatial configuration is such as to lead to the population of this AO equal to 0.33e, the fact being in contradiction with atomic calculations where 6d-state lies much above the ground state. KuCl<sub>4</sub> in comparison with HfCl<sub>4</sub> has a more prominent 7p-character and a less prominent 6d one.

The highest occupied MO (2t<sub>1</sub>) in all these systems is responsible for the ionization potential. These values for KuCl<sub>4</sub> and HfCl<sub>4</sub> will be similar owing to virtually the same value of negative charge on Cl ions and similar 2t<sub>1</sub>-eigenvalues.

By the character of chemical bond KuCl<sub>4</sub> is likely to be close to HfCl<sub>4</sub> than to PbCl<sub>4</sub>, though with a difference in electron density redistribution between d- and p-orbitals. Estimating the chemical bond strength by such an indirect characteristic as crystal field splitting of d-orbitals one can see that the interaction is stronger in HfCl<sub>4</sub> than in KuCl<sub>4</sub> (energy difference between 2e and 6t<sub>2</sub>-orbitals is larger by 0.15 eV in the first molecule compared to the second one). But notwithstanding the close similarity between HfCl<sub>4</sub> and KuCl<sub>4</sub> there is no full analogy between them, especially in the region of antibonding MOs where strong relativistic stabilization of 7s-AO will result in a difference in the spectroscopic properties of KuCl<sub>4</sub>.

Irrespective of what atomic ground state configuration we started from (7s<sup>2</sup>6d<sup>2</sup>, 7s<sup>2</sup>6d7p or 7s<sup>2</sup>7p<sup>2</sup>), the results were approximately the same as a consequence of the SCF procedure.

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#### SAŽETAK

##### Elektronska struktura superteških elemenata i njihovih spojeva

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Samousklađena relativistička  $X_\alpha$  metoda višestrukog raspršenja primijenjena je na sustav  $KuCl_4$ . Ustanovljeno je da  $KuCl_4$  nije potpuni analog spoja  $HfCl_4$ .