CCA-663

539.19:541 Note

Variational Approach to the Correlation Problem in the Molecules

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Received December 21, 1970

The aim of this work is to apply the variational approach as introduced by Da Providencia¹ to the correlation problem. We are interested not only in the calculation of the ground state energy, but, even more, in the excitation transition of the treated systems.

Da Providencia¹ and $\tilde{C}i\check{z}ek^2$ have written the exact wave function in the form:

$$\Psi = (\exp iS) \cdot \Phi$$

where Φ is some appropriate approximation to Ψ and

$$S = \sum_{mi} c_{mi} a_m^+ a_i + \frac{1}{2} \sum_{mnii} c_{mnij} a_m^+ a_n^+ a_j a_i$$

where C_{mi} , C_{mnii} are determined from

$$\langle \Psi / H / \Psi \rangle = min.$$

In our study Φ was taken as a Slater determinant with molecular orbitals determined with the semiempirical SCF-PPP method. The studied systems are

 $CO(NH_2)_2$, $CS(NH_2)_2$, and $CSe(NH_2)_2$.

The previous calculation³ proved that the single — excited states contributions are negligible ($C_{mi} = 0$). Only C_{mnij} were taken in S and were determined from the inhomogeneous system of equations. The calculated energies of the ground states are in Table I.

In the nonheterogenous atom system the doubly — excited states contribution is the largest one but this is not a priori true for heterogeneous one. This fast explains perhaps why the differences E (variational) — E (HF)

Ground State Energies (eV)

	H—F	this work
CO(NH ₂) ₂ CS(NH ₂) ₂ CSe(NH ₂) ₂		106.814 103.527 101.642

are so small. The variational method is very suitable for the calculation of the excitation transitions. With the transformation $H' = e^{-is}He^{is}$ in

$$\mathbf{E} = \langle \Phi \, | \, \mathbf{e}^{-\mathrm{i}\,\mathrm{S}} \, \mathbf{H} \, \mathbf{e}^{-\mathrm{i}\,\mathrm{S}} \, | \, \Phi \rangle$$

we have tried to eliminate the ground state correlation implicit in H. The excitations can be deduced from the transformed Hamiltonian with the time - dependent variational principle. The system of equation is¹

$$\begin{pmatrix} A & B \\ B^x & A^x \end{pmatrix} \begin{pmatrix} X \\ Y \end{pmatrix} = \epsilon \begin{pmatrix} X \\ Y \end{pmatrix}$$

and corresponds formally to the system of the Random-Phase-Approximation (RPA). The difference between both methods is fundamental. While RPA uses Φ as a ground state wave function, the variational method takes (exp S). Φ . The results of both methods are presented in Table II. Both methods

TABLE II

Transition Energies (eV)	Transition Energies (eV)
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	$\Phi \cdot (\partial \mathbf{RPA}) = \Psi$	this work
CO(NH ₂) ₂	of nois	me ap <mark>03.3</mark> priate
$CS(NH_2)_2$	4.93	4.92
CSe(NH ₂) ₂	4.07	4.07

give almost identical results. This is another consequence of the fact that doubly - excited states (in part also quadruply excited) are not involved in Ψ to any large extent. Determining determined as a cost as Ψ yours not all

determined with the semiempicies are reflect. The studied systems

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2. J. Čižek, J. Chem. Phys. 45 (1966) 4256.

3. A. Ažman, M. Drofenik, D. Hadži, and B. Lukman, J. Mol. Structure **1** (1967) 181.

the ground states are in Table I. -mon zeta Variacijska metoda določevanja korelacijske energije molekul ta t

tribution is the largest one **AžMa**. A in **B. Lukman in A. Ažman** one beterogeneous one

Variacijsko metodo¹ smo uporabili za določitev energije osnovnega stanja molekul. Metoda je zelo primerna za račun eksitacij, če je doprinos dvojno vzbujenih stanj k osnovnemu stanju znaten.

Sprejeto 21. prosinca 1970

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