

q -FRACTIONAL DIRAC TYPE SYSTEMS

BILENDER P. ALLAHVERDIEV AND HÜSEYİN TUNA

ABSTRACT. This paper is devoted to study a regular q -fractional Dirac type system. We investigate the properties of the eigenvalues and the eigenfunctions of this system. By using a fixed point theorem we give a sufficient condition on eigenvalues for the existence and uniqueness of the associated eigenfunctions.

1. INTRODUCTION

q -calculus deals with the investigation and applications of quantum derivatives and quantum integrals. It is an interesting topic having interconnections with various problems of mathematical physics and quantum mechanics ([8–10, 12, 14, 16, 17, 23, 24, 30]). For the q -calculus, we refer the reader to the books [7, 13, 18].

The fractional q -calculus is the generalization of the q -calculus. In the recent years, some results have been derived in q -fractional equations [5–7, 15, 20–22, 25, 26]. Mansour [25] introduced q -fractional Sturm-Liouville problems containing the left-sided Caputo q -fractional derivative and the right-sided Riemann-Liouville q -fractional derivative. The author used a fixed point theorem to introduce a sufficient condition on eigenvalues for the existence and uniqueness of the associated eigenfunctions of q -fractional Sturm-Liouville problems. Al-Towailb studied the regular q -fractional Sturm-Liouville problems. The author proved properties of the eigenvalues and the eigenfunctions in [5]. In [26], the author introduced the essential q -fractional variational analysis needed in proving the existence of a countable set of real eigenvalues and associated orthogonal eigenfunctions for the regular q -fractional Sturm-Liouville problems. Allahverdiev and Tuna [3] proved a theorem on the completeness of the system of eigenvectors and associated vectors of the dissipative q -fractional Sturm-Liouville operators.

2020 *Mathematics Subject Classification.* 39A13, 34L40, 34L10, 26A33.

Key words and phrases. q -fractional Dirac operator, eigenvalues, eigenfunctions.

It is well known that the Dirac systems defined by

$$(1.1) \quad \begin{pmatrix} 0 & -\frac{d}{dx} \\ \frac{d}{dx} & 0 \end{pmatrix} \begin{pmatrix} y_1(x) \\ y_2(x) \end{pmatrix} + \begin{pmatrix} p(x) & 0 \\ 0 & r(x) \end{pmatrix} \begin{pmatrix} y_1(x) \\ y_2(x) \end{pmatrix} \\ = \lambda \begin{pmatrix} y_1(x) \\ y_2(x) \end{pmatrix},$$

where $x \in [a, b]$, play an important role in relativistic quantum mechanics. These systems describe spin 1/2 particles, including electrons, neutrinos, muons, protons, neutrons, quarks, and their corresponding anti-particles. For the history and details of the Dirac systems, see [11, 19, 29, 31] and their references. In this paper, we are interested in a q -fractional version of the system (1.1) defined by

$$\begin{pmatrix} 0 & -\mathcal{D}_{q,a^-}^\alpha \\ {}^c\mathcal{D}_{q,0^+}^\alpha & 0 \end{pmatrix} \begin{pmatrix} y_1(x) \\ y_2(x) \end{pmatrix} + \begin{pmatrix} p(x) & 0 \\ 0 & r(x) \end{pmatrix} \begin{pmatrix} y_1(x) \\ y_2(x) \end{pmatrix} \\ = \lambda \begin{pmatrix} \omega_1(x) & 0 \\ 0 & \omega_2(x) \end{pmatrix} \begin{pmatrix} y_1(x) \\ y_2(x) \end{pmatrix},$$

where $x \in (0, a)$. To the best of the authors' knowledge there are no results available in the literature considering this system. These results are a generalization of the regular q -Dirac system introduced in [2].

2. PRELIMINARIES

First of all, we recall the notations and some basic properties for q -fractional calculus theory, which are useful in the following discussion (see [1, 4, 7, 13, 18, 25, 27, 28]). Throughout this paper, we assume that $0 < q < 1$ and A is a q -geometric set, i.e., $qx \in A$ whenever $x \in A$. For every $t > 0$, we define the sets $A_{t,q}$, $A_{t,q}^*$ and $\mathcal{A}_{t,q}$, respectively, by

$$A_{t,q} := \{tq^n : n \in \mathbb{N}\}, \quad A_{t,q}^* := A_{t,q} \cup \{0\},$$

and

$$\mathcal{A}_{t,q} := \{\pm tq^n : n \in \mathbb{N}\}.$$

Let $y(\cdot)$ be a complex-valued function on A . The q -difference operator \mathcal{D}_q is defined by

$$\mathcal{D}_q y(x) = \frac{y(qx) - y(x)}{(q-1)x} \quad \text{for all } x \in A \setminus \{0\}.$$

The q -derivative at zero is defined by

$$\mathcal{D}_q y(0) = \lim_{n \rightarrow \infty} \frac{y(q^n x) - y(0)}{q^n x} \quad (x \in A),$$

if the limit exists and does not depend on x . A *right-inverse* to D_q , the *Jackson q -integration* is given by

$$\int_0^x f(t) d_q t = x(1-q) \sum_{n=0}^{\infty} q^n f(q^n x) \quad (x \in A),$$

provided that the series converges, and

$$\int_a^b f(t) d_q t = \int_0^b f(t) d_q t - \int_0^a f(t) d_q t \quad (a, b \in A).$$

Let $L_q^2(0, a)$ be the space of all complex-valued functions defined on $[0, a]$ such that

$$\|f\| := \left(\int_0^a |f(x)|^2 d_q x \right)^{1/2} < \infty.$$

$L_q^2(0, a)$ is a separable Hilbert space with the inner product

$$(f, g) := \int_0^a f(x) \overline{g(x)} d_q x, \quad f, g \in L_q^2(0, a),$$

and the orthonormal basis

$$\phi_n(x) = \begin{cases} \frac{1}{\sqrt{x(1-q)}}, & x = aq^n, \\ 0, & \text{otherwise,} \end{cases}$$

where $n = 0, 1, 2, \dots$ (see [7]).

DEFINITION 2.1. A function f which is defined on A , $0 \in A$, is said to be q -regular at zero if

$$\lim_{n \rightarrow \infty} f(xq^n) = f(0)$$

for every $x \in A$ (see [7]).

Let $C(A)$ denote the space of all q -regular at zero functions on A . This space is a normed space with the norm function

$$\|f\| = \sup \{|f(xq^n)|, x \in A, n \in \mathbb{N}\}.$$

(see [7]).

DEFINITION 2.2. A q -regular at zero function f which is defined on $A_{t,q}^*$ is said to be q -absolutely continuous if

$$\sum_{j=0}^{\infty} |f(uq^j) - f(uq^{j+1})| \leq K, \quad \forall u \in A_{t,q}^*,$$

for K is a constant depending on the function f (see [7]).

The space of all q -absolutely continuous functions on $A_{t,q}^*$ is denoted by $AC_q(A_{t,q}^*)$. Note that $AC_q(A_{q,t}^*) \subseteq C(A_{q,t}^*)$.

For $n \in \mathbb{N}$ and $\alpha, a_1, \dots, a_n \in \mathbb{C}$; the q -shifted factorial, the multiple q -shifted factorial and the q -binomial coefficients are defined by

$$(a; q)_0 = 1, \quad (a; q)_n = \prod_{k=0}^{n-1} (1 - aq^k), \quad (a; q)_\infty = \prod_{k=0}^{\infty} (1 - aq^k),$$

$$(a_1, a_2, \dots, a_k : q) = \prod_{j=1}^k (a_j; q)_n$$

and

$$\begin{bmatrix} \alpha \\ 0 \end{bmatrix}_q = 1, \quad \begin{bmatrix} \alpha \\ n \end{bmatrix}_q = \frac{(1 - q^\alpha)(1 - q^{\alpha-1}) \dots (1 - q^{\alpha-n+1})}{(q; q)_n},$$

respectively (see [7]). The *generalized q -shifted factorial* is defined by

$$(a; q)_\nu = \frac{(a; q)_\infty}{(aq^\nu; q)_\infty} \quad (\nu \in \mathbb{R})$$

(see [7]). The *q -Gamma function* is defined by

$$\Gamma_q(z) = \frac{(q; q)_\infty}{(q^z; q)_\infty} (1 - q)^{1-z}, \quad z \in \mathbb{C}, \quad |q| < 1$$

(see [7]).

DEFINITION 2.3. *Let $0 < \alpha \leq 1$. The left-sided and right-sided Riemann-Liouville q -fractional operator are given by the formulas*

$$(2.1) \quad \mathcal{I}_{q,a^+}^\alpha f(x) = \frac{x^{\alpha-1}}{\Gamma_q(\alpha)} \int_a^x \left(\frac{qt}{x}; q \right)_{\alpha-1} f(t) d_q t,$$

$$(2.2) \quad \mathcal{I}_{q,b^-}^\alpha f(x) = \frac{1}{\Gamma_q(\alpha)} \int_{qx}^b t^{\alpha-1} \left(\frac{qx}{t}; q \right)_{\alpha-1} f(t) d_q t,$$

respectively (see [25]).

DEFINITION 2.4. *Let $\alpha > 0$ and $[\alpha] = m$. The left-sided and right-sided Riemann-Liouville fractional q -derivatives of the order α are defined, respectively, as follows:*

$$(2.3) \quad \mathcal{D}_{q,a^+}^\alpha f(x) = \mathcal{D}_q^m \mathcal{I}_{q,a^+}^{m-\alpha} f(x),$$

$$(2.4) \quad \mathcal{D}_{q,b^-}^\alpha f(x) = \left(\frac{-1}{q} \right)^m \mathcal{D}_{q^{-1}}^m \mathcal{I}_{q,b^-}^{m-\alpha} f(x).$$

Similar formulas give the left-sided and right-sided Caputo fractional q -derivatives of order α , respectively as follows:

$$\begin{aligned} {}^c\mathcal{D}_{q,a^+}^\alpha f(x) &= \mathcal{I}_{q,a^+}^{m-\alpha} \mathcal{D}_q^m f(x), \\ {}^c\mathcal{D}_{q,b^-}^\alpha f(x) &= \left(\frac{-1}{q}\right)^m \mathcal{I}_{q,b^-}^{m-\alpha} \mathcal{D}_{q^{-1}}^m f(x) \end{aligned}$$

(see [25]).

In order to prove the main results, we also need the following lemmas. One can find them in [25].

LEMMA 2.5.

i) The left-sided Riemann-Liouville q -fractional operator satisfies the semi-group property

$$(2.5) \quad \mathcal{I}_{q,a^+}^\alpha \mathcal{I}_{q,a^+}^\beta = \mathcal{I}_{q,a^+}^{\alpha+\beta} f(x), \quad x \in A_{q,a}^*,$$

for any function defined on $A_{q,a}$ and for any values of α and β .

ii) The right-sided Riemann-Liouville q -fractional operator satisfies the semi-group property

$$\mathcal{I}_{q,b^-}^\alpha \mathcal{I}_{q,b^-}^\beta f(x) = \mathcal{I}_{q,b^-}^{\alpha+\beta} f(x), \quad x \in A_{q,b}^*,$$

for any function defined on $A_{q,b}$ and for any values of α and β .

LEMMA 2.6. Let $\alpha \in (0, 1)$.

i) If $f \in L_q^1(A_{q,a}^*)$ such that $\mathcal{I}_{q,0^+}^\alpha f \in AC_q(A_{t,q}^*)$ then

$${}^c\mathcal{D}_{q,0^+}^\alpha \mathcal{I}_{q,0^+}^\alpha f(x) = f(x) - \frac{\mathcal{I}_{q,0^+}^\alpha f(0)}{\Gamma_q(1-\alpha)} x^{-\alpha}.$$

Moreover, if f is bounded on $A_{t,q}^*$ then

$${}^c\mathcal{D}_{q,0^+}^\alpha \mathcal{I}_{q,0^+}^\alpha f(x) = f(x).$$

ii) If $f \in L_q^1(A_{q,a})$ then

$$\mathcal{D}_{q,0^+}^\alpha \mathcal{I}_{q,0^+}^\alpha f(x) = f(x).$$

iii) If f is a function defined on $A_{t,q}^*$ then

$${}^c\mathcal{D}_{q,a^-}^\alpha \mathcal{I}_{q,a^-}^\alpha f(x) = f(x) - \frac{a^{-\alpha}}{\Gamma_q(1-\alpha)} \left(\frac{qx}{a}; q\right)_{-\alpha} \left(\mathcal{I}_{q,a^-}^\alpha f\right) \left(\frac{a}{q}\right),$$

$$\mathcal{D}_{q,a^-}^\alpha \mathcal{I}_{q,a^-}^\alpha f(x) = f(x),$$

$$\mathcal{I}_{q,a^-}^\alpha \mathcal{D}_{q,a^-}^\alpha f(x) = f(x) - \frac{a^{\alpha-1}}{\Gamma_q(\alpha)} \left(\frac{qx}{a}; q\right)_{\alpha-1} \left(\mathcal{I}_{q,a^-}^{1-\alpha} f\right) \left(\frac{a}{q}\right).$$

iv) If $f \in AC_q(A_{t,q}^*)$ then

$$\mathcal{I}_{q,0^+}^\alpha {}^c\mathcal{D}_{q,0^+}^\alpha f(x) = f(x) - f(0).$$

We denote by $L_{q,\omega}^2(A_{t,\alpha}^*; E)$ ($E := \mathbb{R}^2$) the Hilbert space which consists of vector-valued functions with inner product

$$(2.6) \quad (f, g) := \int_0^a f_1(x) \overline{g_1(x)} \omega_1(x) d_q x + \int_0^a f_2(x) \overline{g_2(x)} \omega_2(x) d_q x,$$

where $f(x) = \begin{pmatrix} f_1(x) \\ f_2(x) \end{pmatrix}$, $g(x) = \begin{pmatrix} g_1(x) \\ g_2(x) \end{pmatrix}$, $f_i(x)$, $g_i(x)$, $\omega_{i\alpha}(x)$ ($i = 1, 2$) are real-valued functions on $A_{t,\alpha}^*$ and $\omega_i(x) > 0$, $\forall x \in A_{t,\alpha}^*$, ($i = 1, 2$).

3. q -FRACTIONAL DIRAC SYSTEMS

In the present section, our goal is to study the q -fractional Dirac system which includes the right-sided Caputo and the left-sided Riemann-Liouville fractional derivatives of same order α . Throughout this section, we assume $\alpha \in (0, 1)$.

Let

$$\begin{aligned} \tau_{q,\alpha} y &:= \begin{pmatrix} 0 & -\mathcal{D}_{q,a^-}^\alpha \\ {}^c\mathcal{D}_{q,0^+}^\alpha & 0 \end{pmatrix} \begin{pmatrix} y_1 \\ y_2 \end{pmatrix} + \begin{pmatrix} p(x) & 0 \\ 0 & r(x) \end{pmatrix} \begin{pmatrix} y_1 \\ y_2 \end{pmatrix} \\ &= \begin{pmatrix} -\mathcal{D}_{q,a^-}^\alpha y_2 + p(x) y_1 \\ {}^c\mathcal{D}_{q,0^+}^\alpha y_1 + r(x) y_2 \end{pmatrix}, \end{aligned}$$

where $y := \begin{pmatrix} y_1 \\ y_2 \end{pmatrix}$. With this notation, we consider the q -fractional Dirac type system:

$$(3.1) \quad \tau_{q,\alpha} f_\lambda = \lambda \omega f_\lambda, \quad a \leq x \leq b < \infty,$$

where $f_\lambda = \begin{pmatrix} f_{\lambda 1} \\ f_{\lambda 2} \end{pmatrix}$, $p(\cdot)$, $r(\cdot)$ are real-valued functions defined in $A_{t,\alpha}^*$, $\omega(x) = \begin{pmatrix} \omega_1(x) & 0 \\ 0 & \omega_2(x) \end{pmatrix}$, $\omega_i(\cdot)$ are real-valued functions defined in $A_{t,\alpha}^*$ and $\omega_{i\alpha}(x) > 0$, $\forall x \in A_{t,\alpha}^*$, ($i = 1, 2$), λ is a complex eigenvalue parameter and boundary conditions

$$(3.2) \quad c_{11} f_{\lambda 1}(0) + c_{12} \mathcal{I}_{q,a^-}^{1-\alpha} f_{\lambda 2}(0) = 0,$$

$$(3.3) \quad c_{21} f_{\lambda 1}(a) + c_{22} \mathcal{I}_{q,a^-}^{1-\alpha} f_{\lambda 2}\left(\frac{a}{q}\right) = 0,$$

with $c_{11}^2 + c_{12}^2 \neq 0$ and $c_{21}^2 + c_{22}^2 \neq 0$.

To pass from the differential expression $T_{q,\alpha} := \omega^{-1} \tau_{q,\alpha}$ to operators, we introduce the space $H \subseteq L_{q,\omega}^2(A_{t,\alpha}^*; E) \cap C(A_{t,\alpha}^*; E)$ which consists of all q -regular at zero functions satisfying the conditions (3.2) and (3.3) with inner product (2.6).

THEOREM 3.1. *The operator $T_{q,\alpha}$ generated by q -fractional Dirac type system (FD) defined by (3.1)-(3.3) is formally self-adjoint on H .*

PROOF. Let $u(\cdot), z(\cdot) \in H$. Then, we have

$$\begin{aligned} (T_{q,\alpha}u, z) - (u, T_{q,\alpha}z) &= \int_0^a ({}^c\mathcal{D}_{q,0^+}^\alpha u_1 + r(x)u_2) \overline{z_2} d_q x \\ &\quad + \int_0^a (-\mathcal{D}_{q,a^-}^\alpha u_2 + p(x)u_1) \overline{z_1} d_q x \\ &\quad - \int_0^a u_2 \overline{({}^c\mathcal{D}_{q,0^+}^\alpha z_1 + r(x)z_2)} d_q x \\ &\quad - \int_0^a u_1 \overline{(-\mathcal{D}_{q,a^-}^\alpha z_2 + p(x)z_1)} d_q x \\ &= \int_0^a ({}^c\mathcal{D}_{q,0^+}^\alpha u_1) \overline{z_2} d_q x - \int_0^a (\mathcal{D}_{q,a^-}^\alpha u_2) \overline{z_1} d_q x \\ &\quad - \int_0^a u_2 \overline{({}^c\mathcal{D}_{q,0^+}^\alpha z_1)} d_q x + \int_0^a u_1 \overline{(\mathcal{D}_{q,a^-}^\alpha z_2)} d_q x. \end{aligned}$$

Since

$$\begin{aligned} \int_0^a ({}^c\mathcal{D}_{q,0^+}^\alpha u_1) \overline{z_2} d_q x &= \int_0^a u_1 \overline{(-\mathcal{D}_{q,a^-}^\alpha z_1)} d_q x \\ &- \left[u_1(a) \overline{\mathcal{I}_{q,a^-}^{1-\alpha} z_2 \left(\frac{a}{q}\right)} - u_1(0) \overline{\mathcal{I}_{q,a^-}^{1-\alpha} z_2(0)} \right] \end{aligned}$$

and

$$\begin{aligned} \int_0^a u_2 \overline{({}^c\mathcal{D}_{q,0^+}^\alpha z_1)} d_q x &= \int_0^a (-\mathcal{D}_{q,a^-}^\alpha u_2) \overline{z_1} d_q x \\ &- \left[z_1(a) \overline{\mathcal{I}_{q,a^-}^{1-\alpha} u_2 \left(\frac{a}{q}\right)} - z_1(0) \overline{\mathcal{I}_{q,a^-}^{1-\alpha} u_2(0)} \right], \end{aligned}$$

we get

$$(3.4) \quad (T_{q,\alpha}u, z) - (u, T_{q,\alpha}z) = [u, z](a) - [u, z](0),$$

where $[y, z](x) := y_1(x) \overline{\mathcal{I}_{q,a^-}^{1-\alpha} z_2(x)} - \overline{z_1(x)} \mathcal{I}_{q,a^-}^{1-\alpha} y_2(x)$. We proceed to show that the equality $(T_{q,\alpha}u, z) = (u, T_{q,\alpha}z)$ for any $u(\cdot), z(\cdot) \in H$. From the boundary conditions (3.2) and (3.3), we get $[u, z]_a = 0$ and $[u, z]_0 = 0$. Consequently,

$$(3.5) \quad (T_{q,\alpha}u, z) = (u, T_{q,\alpha}z).$$

This completes the proof. \square

LEMMA 3.2. *All eigenvalues of the operator $T_{q,\alpha}$ generated by q -FD system defined by (3.1)-(3.3) are real.*

PROOF. Let μ be an eigenvalue with an eigenfunction $z(x)$. From the equality (3.5), we get

$$(3.6) \quad (T_{q,\alpha}z, z) = (z, T_{q,\alpha}z) = (z, \mu z) = \bar{\mu}(z, z).$$

On the other hand,

$$(3.7) \quad (T_{q,\alpha}z, z) = (\mu z, z) = \mu(z, z).$$

It follows from (3.6) and (3.7) that

$$\mu(z, z) = \bar{\mu}(z, z), \quad (\mu - \bar{\mu})(z, z) = 0.$$

Since $z \neq 0$, we get $\mu = \bar{\mu}$. \square

LEMMA 3.3. *If μ_1 and μ_2 are two different eigenvalues of the operator $T_{q,\alpha}$ generated by q -FD system defined by (3.1)-(3.3), then the corresponding eigenfunctions θ and η are orthogonal.*

PROOF. Let μ_1 and μ_2 be two different real eigenvalues with corresponding eigenfunctions θ and η , respectively. From (3.5), we obtain

$$(T_{q,\alpha}\theta, \eta) = (\theta, T_{q,\alpha}\eta), \quad (\mu_1\theta, \eta) = (\theta, \mu_2\eta)$$

$$(\mu_1 - \mu_2)(\theta, \eta) = 0.$$

Since $\mu_1 \neq \mu_2$, we obtain that $\theta(x)$ and $\eta(x)$ are orthogonal. \square

Now let $u(x) = \begin{pmatrix} u_1(x) \\ u_2(x) \end{pmatrix}$, $z(x) = \begin{pmatrix} z_1(x) \\ z_2(x) \end{pmatrix} \in H$. Then, we define the Wronskian of $u(x)$ and $z(x)$ by

$$W(u, z)(x) = u_1(x) \mathcal{I}_{q,a}^{1-\alpha} z_2(x) - z_1(x) \mathcal{I}_{q,a}^{1-\alpha} u_2(x).$$

THEOREM 3.4. *The Wronskian of any solution of Eq. (3.1) is independent of x .*

PROOF. Let $u(x)$ and $z(x)$ be two solutions of Eq. (3.1). By Green's formula (3.4), we have

$$(T_{q,\alpha}u, z) - (u, T_{q,\alpha}z) = [u, z](a) - [u, z](0).$$

Since $T_{q,\alpha}u = \lambda u$ and $T_{q,\alpha}z = \lambda z$, we have

$$(\lambda u, z) - (u, \lambda z) = [u, z](a) - [u, z](0),$$

$$(\lambda - \bar{\lambda})(u, z) = [u, z](a) - [u, z](0).$$

Since $\lambda \in \mathbb{R}$, we have $[u, z](a) = [u, z](0) = W(u, \bar{z})(0)$, i.e., the Wronskian is independent of x . \square

COROLLARY 3.5. *If $u(x)$ and $z(x)$ are both solutions of Equation (3.1), then either $W(u, z)(x) = 0$ or $W(u, z)(x) \neq 0$ for all $x \in [0, a]$.*

THEOREM 3.6. *Any two solutions of Equation (3.1) are linearly dependent if and only if their Wronskian is zero.*

PROOF. Let $u(x)$ and $z(x)$ be two linearly dependent solutions of equation (3.1). Then, there exists a constant $k > 0$ such that $u(x) = k z(x)$. Hence

$$W(u, z) = \begin{vmatrix} u_1(x) & \mathcal{I}_{q,a^-}^{1-\alpha} u_2(x) \\ z_1(x) & \mathcal{I}_{q,a^-}^{1-\alpha} z_2(x) \end{vmatrix} = \begin{vmatrix} kz_1(x) & k\mathcal{I}_{q,a^-}^{1-\alpha} z_2(x) \\ z_1(x) & \mathcal{I}_{q,a^-}^{1-\alpha} z_2(x) \end{vmatrix} = 0.$$

Conversely, the Wronskian $W(u, z) = 0$ and therefore, $u(x) = kz(x)$, i.e., $u(x)$ and $z(x)$ are linearly dependent. \square

Before proceeding further, we need the following auxiliary functions.

We introduce the function $\phi(x) := \begin{pmatrix} (\mathcal{I}_{q,a^-}^\alpha - 1)(x) \\ (\mathcal{I}_{q,0^+}^\alpha + 1)(x) \end{pmatrix}$. Further, the general solution of the equation $\tau_{q,\alpha}\psi = 0$, i.e.,

$$\begin{pmatrix} 0 & \mathcal{D}_{q,a^-}^\alpha \\ {}^c\mathcal{D}_{q,0^+}^\alpha & 0 \end{pmatrix} \begin{pmatrix} \psi_1 \\ \psi_2 \end{pmatrix} = 0$$

is given by

$$\psi = \begin{pmatrix} \xi_1 \\ \xi_2 \varphi(\alpha, a, x) \end{pmatrix},$$

where

$$(3.8) \quad \varphi(\alpha, a, x) = \frac{a^{\alpha-1} \left(\frac{qx}{a} : q\right)_{\alpha-1}}{\Gamma_q(\alpha)}.$$

LEMMA 3.7. *Let*

$$\Delta := c_{11}c_{12} - c_{11}c_{21}$$

and

$$(3.9) \quad F_\lambda(f) := \{V - \lambda\omega\} f_\lambda,$$

where $V(x) := \begin{pmatrix} p(x) & 0 \\ 0 & r(x) \end{pmatrix}$. Assume $\Delta \neq 0$. Then on the space $C(A_{t,\alpha}^*)$, the q -FD system defined by (3.1)-(3.3) is equivalent to the integral equation

$$f_\lambda(x) = -MF_\lambda(f) + A(x)T + B(x)Z,$$

where the coefficients M, A, T, B and Z are

$$M := \begin{pmatrix} 0 & \mathcal{I}_{q,0^+}^\alpha \\ \mathcal{I}_{q,a^-}^\alpha & 0 \end{pmatrix},$$

$$A(x) := \begin{pmatrix} \frac{c_{12}c_{22}}{\Delta} \\ -\frac{c_{21}c_{12}}{\Delta}\varphi(\alpha, a, x) \end{pmatrix},$$

$$T := -\mathcal{I}_{q,a^-}^\alpha F_{\lambda 1}(y) |_{x=0},$$

$$B(x) := \begin{pmatrix} \frac{c_{12}c_{21}}{\Delta} \\ -\frac{c_{21}c_{11}}{\Delta}\varphi(\alpha, a, x) \end{pmatrix},$$

$$Z := -\mathcal{I}_{q,0^+}^1 F_{\lambda 2}(y) |_{x=a},$$

and the function $\varphi(\alpha, a, x)$ is defined in (3.8).

PROOF. Using fractional composition rules and (3.9), we can rewrite the equation (3.1) as follows:

$$\tau_{q,\alpha} [f_\lambda(x) + MF_\lambda(f)] = 0.$$

Thus, we get

$$f_\lambda(x) + MF_\lambda(f) = \begin{pmatrix} \xi_1 \\ \xi_2\varphi(\alpha, a, x) \end{pmatrix},$$

i.e.,

$$(3.10) \quad f_\lambda(x) = -MF_\lambda(f) + \begin{pmatrix} \xi_1 \\ \xi_2\varphi(\alpha, a, x) \end{pmatrix}.$$

Now, we shall connect the coefficients ξ_i ($i = 1, 2$) to the values c_{ij} ($i, j = 1, 2$) in the boundary conditions (3.2)-(3.3). From the equation (3.10), we obtain

$$Kf_\lambda(x) = -KMf_\lambda(f) + K \begin{pmatrix} \xi_1 \\ \xi_2\varphi(\alpha, a, x) \end{pmatrix},$$

where $K := \begin{pmatrix} 0 & \mathcal{I}_{q,a^-}^{1-\alpha} \\ 1 & 0 \end{pmatrix}$. Then we have

$$\begin{pmatrix} \mathcal{I}_{q,a^-}^{1-\alpha} f_{\lambda 2} \\ f_{\lambda 1} \end{pmatrix} = - \begin{pmatrix} \mathcal{I}_{q,a^-}^1 & 0 \\ 0 & \mathcal{I}_{q,0^+}^\alpha \end{pmatrix} F_\lambda(f) + \begin{pmatrix} \mathcal{I}_{q,a^-}^{1-\alpha} [\xi_2\varphi(\alpha, a, x)] \\ \xi_1 \end{pmatrix},$$

i.e.,

$$\begin{pmatrix} \mathcal{I}_{q,a^-}^{1-\alpha} f_{\lambda 2} \\ f_{\lambda 1} \end{pmatrix} = \begin{pmatrix} -\mathcal{I}_{q,a^-}^1 F_{\lambda 1}(f) \\ -\mathcal{I}_{q,0^+}^\alpha F_{\lambda 2}(f) \end{pmatrix} + \begin{pmatrix} \xi_2 \\ \xi_1 \end{pmatrix}.$$

By virtue of (3.2) and (3.3), we conclude that

$$\begin{aligned} f_{\lambda 1}(0) &= \xi_1, \\ f_{\lambda 1}(a) &= -\mathcal{I}_{q,0^+}^\alpha F_{\lambda 2}(y)|_{x=a} + \xi_1, \\ \mathcal{I}_{q,a^-}^{1-\alpha} f_{\lambda 2}(0) &= -\mathcal{I}_{q,a^-}^1 F_{\lambda 1}(y)|_{x=0} + \xi_2, \\ \mathcal{I}_{q,a^-}^{1-\alpha} f_{\lambda 2}\left(\frac{a}{q}\right) &= \xi_2. \end{aligned}$$

This leads to the system of equations

$$c_{11}\xi_1 + c_{12}\xi_2 = -c_{12}, \quad Tc_{21}\xi_1 + c_{22}\xi_2 = -c_{21}Z.$$

Since $\Delta \neq 0$, the solutions for coefficients $\xi_j, j = 1, 2$ is unique:

$$\begin{aligned} \xi_1 &= \frac{c_{12}(c_{21}Z - c_{22}T)}{\Delta}, \\ \xi_2 &= \frac{c_{21}(c_{12}T - c_{11}Z)}{\Delta}. \end{aligned}$$

We have finished the proof of the lemma. □

Now, we prove the existence and uniqueness of eigenfunction of the regular q -FD system defined by (3.1)-(3.3). In the next result, we use the following notations:

$$A := \|A(x)\|_C, \quad B := \|B(x)\|_C, \quad S_\phi := \|\phi(x)\|_C,$$

where $\|\cdot\|_C$ denotes the supremum norm on the space $C(A_{t,\alpha}^*, E)$.

THEOREM 3.8. *Let $\alpha \in (0, 1)$ and assume $\Delta \neq 0$. Then unique continuous function y_λ for the regular q -FD system defined by (3.1)-(3.3) corresponding to each eigenvalue obeying*

$$(3.11) \quad \|V - \lambda\omega\|_C \leq \frac{1}{S_\phi + A\|\phi(a)\|_C + Ba}$$

exists and such eigenvalue is simple.

PROOF. Let us define the mapping $L : C(A_{t,\alpha}^*, E) \rightarrow C(A_{t,\alpha}^*, E)$ by

$$Lf := -MF_\lambda(f) + A(x)T + B(x)Z.$$

Now, we show that the equation (3.1) can be interpreted as a fixed point condition on the space $C(A_{t,\alpha}^*, E)$. Using the following estimate

$$\|F_\lambda(g) - F_\lambda(h)\|_C \leq \|g - h\|_C \|V - \lambda\omega\|_C,$$

we conclude that

$$\begin{aligned} \|Lg - Lh\|_C &\leq \|g - h\|_C \|V - \lambda\omega\|_C S_\phi + A \|g - h\|_C \|\phi(a)\|_C \\ &\quad + Ba \|g - h\|_C \|V - \lambda\omega\|_C \\ &= \|V - \lambda\omega\|_C \|g - h\|_C (S_\phi + A \|\phi(a)\|_C + Ba) \\ &= \Pi \|g - h\|_C, \end{aligned}$$

where $\Pi = \|V - \lambda\omega\|_C (S_\phi + A \|\phi(a)\|_C + Ba)$. By the condition (3.11), the mapping L is a contraction on the space $C(A_{t,\alpha}^*, E)$ so it has a unique fixed point. Therefore, such eigenvalue is simple. \square

CONCLUSION 3.9. In this paper, we study regular q -fractional Dirac systems. In this context, we investigate the properties of the eigenvalues and the eigenfunctions of this system. Finally, we give a sufficient condition on eigenvalues for the existence and uniqueness of the associated eigenfunctions.

REFERENCES

- [1] R. P. Agarwal, *Certain fractional q -integrals and q -derivatives*, Proc. Camb. Phil. Soc. **66** (1969), 365–370.
- [2] B. P. Allahverdiev and H. Tuna, *One-dimensional q -Dirac equation*, Math. Meth. Appl. Sci. **40** (2017), 7287–7306.
- [3] B. P. Allahverdiev and H. Tuna, *Spectral analysis of q -fractional Sturm-Liouville operators*, Electr. J. Differential Equations **2017** (2017), Paper No. 136, 17 pp.
- [4] W. A. Al-Salam, *Some fractional q -integrals and q -derivatives*, Proc. Edinb. Math. Soc. (2) **15** (1966/1967), 135–140.
- [5] M. A. Al-Towailb, *A q -fractional approach to the regular Sturm-Liouville problems*, Electron. J. Differential Equations **2017** (2017), Paper No. 88, 13 pp.
- [6] M. A. Al-Towailb, *The solution of certain triple q -integral equations in fractional q -calculus approach*, Arab. J. Math. Sci. **25** (2019), 17–28.
- [7] M. H. Annaby and Z. S. Mansour, *q -Fractional Calculus and Equations*, Lecture Notes in Mathematics, vol. **2056**, Springer, Heidelberg, 2012.
- [8] G. E. Andrews, R. Askey and R. Roy, *Special Functions*, Cambridge Univ. Press, 1999.
- [9] N. M. Atakishiyev, A. U. Klimyk and K. B. Wolf, *A discrete quantum model of the harmonic oscillator*, J. Phys. A: Math. Theor. **41** (2008), no. 8, 085201, 14 pp.
- [10] L. C. Biedenharn, *The quantum group $SU_q(2)$ and a q -analogue of the boson operators*, J. Phys. A: Math. Gen. **22** (1989), L873–L878.
- [11] P. Collas and D. Klein, *The Dirac Equation in Curved Spacetime, A Guide for Calculations*, SpringerBriefs in Physics Series, Springer, Cham, 2019.
- [12] R. Díaz and E. Pariguán, *An example of Feynman–Jackson integrals*, J. Phys. A: Math. Theor. **40** (2007), 1265–1272.
- [13] T. Ernst, *The History of q -Calculus and a New Method*, U. U. D. M. Report (2000):16, Department of Mathematics, Uppsala University, 2000.
- [14] K. Ey, A. Ruffing and S. Suslov, *Method of separation of the variables for basic analogs of equations of mathematical physics*, Ramanujan J. **13** (2007) 407–447.
- [15] R. A. C. Ferreira, *Nontrivial solutions for fractional q -difference boundary value problems*, Electron. J. Qual. Theory Differ. Equ. **2010** (2010), No. 70, 10 pp.
- [16] R. J. Finkelstein, *q -uncertainty relations*, Internat. J. Modern Phys. A **13** (1998), 1795–1803.

- [17] G. Gasper and M. Rahman, *Basic Hypergeometric Series*, second edition, Cambridge Univ. Press, 2004.
- [18] V. Kac and P. Cheung, *Quantum Calculus*, Springer-Verlag, Berlin-Heidelberg, 2002.
- [19] B. M. Levitan and I. S. Sargsjan, *Introduction to Spectral Theory: Self adjoint Ordinary Differential Operators*, American Mathematical Society, Providence, RI, USA, 1975.
- [20] X. Li, Z. Han, S. Sun and H. Lu, *Boundary value problems for fractional q -difference equations with nonlocal conditions*, *Adv. Differ. Equ.* **2014** (2014), Paper No. 57, 16 pp.
- [21] X. Li and Z. Han, *Boundary value problems of fractional q -difference Schrödinger equations*, *Appl. Math. Lett.* **46** (2015), 100–105.
- [22] S. Liang, *Positive solutions for singular boundary value problem with fractional q -differences*, *Bull. Malays. Math. Sci. Soc.* **38** (2015), 647–666.
- [23] A. Lorek, A. Ruffing and J. Wess, *A q -deformation of the harmonic oscillator*, *Z. Phys. C* **74** (1997), 369–377.
- [24] A. J. Macfarlane, *On q -analogues of the quantum harmonic oscillator and the quantum group $SU(2)_q$* , *J. Phys. A: Math. Gen.* **22** (1989), 4581–4588.
- [25] Z. S. I. Mansour, *On fractional q -Sturm-Liouville problems*, *J. Fixed Point Theory Appl.* **19** (2017), 1591–1612.
- [26] Z. S. I. Mansour, *Variational methods for fractional q -Sturm-Liouville problems*, *Bound. Value Probl.* **2016** (2016), Paper No. 150, 31 pp.
- [27] P. M. Rajković, S. D. Marinković and M. S. Stanković, *A generalization of the concept of q -fractional integrals*, *Acta Math. Sin. (Engl. Ser.)* **25** (2009), 1635–1646.
- [28] M. S. Stanković, P. M. Rajković and S. D. Marinković, *On q -fractional order derivatives of Riemann-Liouville and Caputo type*, (2009), <http://arxiv.org/abs/0909.0387>
- [29] B. Thaller, *The Dirac Equation*, Springer-Verlag, Berlin, 1992.
- [30] M. R. Ubriaco, *Time evolution in quantum mechanics on the quantum line*, *Phys. Lett. A* **163** (1992) 1–4.
- [31] J. Weidmann, *Spectral Theory of Ordinary Differential Operators*, Lecture Notes in Mathematics **1258**, Springer, Berlin, 1987.

q -frakcijski sustavi Diracovog tipa

Bilender P. Allahverdiev i Hüseyin Tuna

SAŽETAK. Ovaj članak je posvećen proučavanju regularnih q -frakcijskih sustava Diracovog tipa. Proučavaju se svojstva svojstvenih vrijednosti i svojstvenih funkcija tih sustava. Korištenjem teorema o fiksnoj točki, daje se dovoljan uvjet na svojstvene vrijednosti za postojanje i jedinstvenost pridruženih svojstvenih funkcija.

Bilender P. Allahverdiev
Department of Mathematics
Faculty of Arts and Sciences,
Süleyman Demirel University
32260 Isparta, Turkey
E-mail: bilenderpasaoglu@sdu.edu.tr

Hüseyin Tuna
Department of Mathematics
Faculty of Arts and Sciences
Mehmet Akif Ersoy University
15030 Burdur, Turkey
E-mail: hustuna@gmail.com

Received: 27.1.2020.

Revised: 11.7.2020.

Accepted: 15.9.2020.