

## Explicit forms for three integrals in Wand et al.\*

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Received September 3, 2020; accepted January 13, 2021

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**Abstract.** We derive explicit forms for the three integrals used in Kim and Wand [3] and Wand, Ormerody, Padoan and Frühwirth [7]. The explicit forms involve known special functions for which in-built routines are available.

**AMS subject classifications:** Primary 60E10; Secondary 33C60, 62G32

**Key words:** Grünwald–Letnikov fractional derivative, Kummer’s confluent hypergeometric function, parabolic cylinder function, Tricomi confluent hypergeometric function

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### 1. Introduction

Article [3] gave an explicit form of expectation propagation for a simple statistical model, while [7] studied mean field variational Bayes for elaborate distributions. Their studies involved the following three integrals. The first integral is

$$\mathcal{A}(p, q, r, s, t, u) = \int_{\mathbb{R}} \frac{x^p \exp\{qx - rx^2\}}{(t + sx + x^2)^u} dx \quad (1)$$

for  $p \geq 0, q, s \in \mathbb{R}, r, u > 0$  and  $s^2 < 4t$ ; see equation (2.1) on page 552 of [3]. The second integral is

$$\mathcal{B}(p, q, r, s, t, u) = \int_{\mathbb{R}} \frac{x^p \exp\{qx - re^x - s e^x / (t + e^x)\}}{(t + e^x)^u} dx \quad (2)$$

for  $p, s \geq 0, q \in \mathbb{R}$  and  $r, t, u > 0$ ; see equation (2.1) on page 552 of [3]. The third integral is

$$\mathcal{I}(p, q, r, s) = \int_{\mathbb{R}} x^p \exp\{qx - rx^2 - s e^{-x}\} dx \quad (3)$$

for  $p \geq 0, q \in \mathbb{R}$  and  $s, r > 0$ ; see page 851 in [7].

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\*The research of Tibor K. Pogány has been supported in part by the University of Rijeka, Croatia; project codes `uniri-pr-prirod-19-16` and `uniri-tehnic-18-66`.

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Both publications [3] and [7] provided explicit forms for none of the three integrals. In this note, we provide explicit forms for all three integrals. The forms involve known special functions and in-built routines for computing them are available in the literature.

The organization of this note is the following. Section 2 gives two explicit forms for (1). Section 3 gives one explicit form for (2). Section 4 gives two explicit forms for (3).

## 2. The integral $\mathcal{A}$

We present an approach in which we use the Grünwald-Letnikov fractional derivative. The Grünwald-Letnikov fractional derivative of order  $\nu$  with respect to the argument  $x$  of a suitable function  $f$  is defined by [6]

$$\mathbb{D}_x^\nu[f] := \lim_{h \downarrow 0} \frac{1}{h^\nu} \sum_{m=0}^{\infty} (-1)^m \binom{\nu}{m} f(x + (\nu - m)h),$$

where  $h \downarrow 0$  means that in approaching zero  $h$  remains positive. As is well-known (see, for example, [4]), the Grünwald-Letnikov fractional derivative  $\mathbb{D}_x^\nu$  of order  $\nu$  of the exponential function is

$$\mathbb{D}_x^\nu [e^{\alpha x}] = \alpha^\nu e^{\alpha x}. \quad (4)$$

Firstly, consider the well-known integral

$$\mathcal{I}(\alpha, \beta) = \int_{\mathbb{R}} e^{\alpha x - \beta x^2} dx = \sqrt{\frac{\pi}{\beta}} \exp\left\{\frac{\alpha^2}{4\beta}\right\}, \quad \Re(\beta) > 0.$$

Obviously, we have

$$\mathbb{D}_\alpha^p [\mathcal{I}(\alpha, \beta)] = \int_{\mathbb{R}} x^p e^{\alpha x - \beta x^2} dx.$$

On the other hand,

$$\mathbb{D}_\alpha^p [\mathcal{I}(\alpha, \beta)] = \int_0^\infty x^p e^{\alpha x - \beta x^2} dx + \int_{-\infty}^0 x^p e^{\alpha x - \beta x^2} dx =: I^+ + I^-.$$

Using equation (13), page 313 of [1],

$$I^+ = (2\beta)^{-\frac{p+1}{2}} \Gamma(p+1) \exp\left\{\frac{\alpha^2}{8\beta}\right\} D_{-p-1}\left(-\frac{\alpha}{\sqrt{2\beta}}\right),$$

where  $D_\mu(\cdot)$  denotes the parabolic cylinder function of order  $\mu$  (see, for example, [2]), and the constraint  $p+1 > 0$  should be satisfied (which is definitely a weaker assumption than the assumed  $p > 0$  by Kim and Wand [3]). Accordingly, we have

$$I^- = (-1)^p (2\beta)^{-\frac{p+1}{2}} \Gamma(p+1) \exp\left\{\frac{\alpha^2}{8\beta}\right\} D_{-p-1}\left(\frac{\alpha}{\sqrt{2\beta}}\right).$$

Therefore,

$$\begin{aligned} \mathbb{D}_\alpha^p [\mathcal{I}(\alpha, \beta)] &= (2\beta)^{-\frac{p+1}{2}} \Gamma(p+1) \exp\left\{\frac{\alpha^2}{8\beta}\right\} \\ &\quad \times \left[ D_{-p-1}\left(-\frac{\alpha}{\sqrt{2\beta}}\right) + (-1)^p D_{-p-1}\left(\frac{\alpha}{\sqrt{2\beta}}\right) \right]. \end{aligned}$$

Now, introduce a parameter  $a > 0$  and specify  $\alpha = q - as \in \mathbb{R}$  and  $\beta = r + a$ , the latter evidently positive; therefore nothing harms the assumption on the parameter space of  $(p, q, r, s, t)$ . Considering now the integral

$$e^{-at} \mathbb{D}_{q-as}^p [\mathcal{I}(q - as, r + a)] = \int_{\mathbb{R}} x^p e^{qx - rx^2 - a(t+sx+x^2)} dx,$$

we conclude that

$$\mathcal{A}(p, q, r, s, t, u) = (-1)^u \mathbb{D}_a^{-u} [e^{-at} \mathbb{D}_{q-as}^p [\mathcal{I}(q - as, r + a)]] \Big|_{a=0}.$$

This formula proves the following result.

**Proposition 1.** *For all  $p \geq 0, q, s \in \mathbb{R}, r, u > 0$  and  $s^2 < 4t$ , we have*

$$\begin{aligned} \mathcal{A}(p, q, r, s, t, u) &= e^{i\pi u} 2^{-\frac{p+1}{2}} \Gamma(p+1) \lim_{a \downarrow 0} \mathbb{D}_a^{-u} \left[ (r+a)^{-\frac{p+1}{2}} \exp\left\{\frac{(q-as)^2}{8(r+a)} - at\right\} \right. \\ &\quad \left. \times \left\{ D_{-p-1}\left(\frac{as-q}{\sqrt{2(r+a)}}\right) + e^{i\pi p} D_{-p-1}\left(\frac{q-as}{\sqrt{2(r+a)}}\right) \right\} \right]. \end{aligned}$$

We now present another approach to calculating integral  $\mathcal{A}$ . Kummer's (or confluent hypergeometric) function series definition is

$${}_1F_1(a, c, z) = \sum_{n \geq 0} \frac{(a)_n}{(c)_n} \frac{z^n}{n!}.$$

The parabolic cylinder function  $D_\nu$  is expressible in terms of the Tricomi confluent hypergeometric function, *viz.*

$$U(a; c; z) = \frac{\Gamma(1-c)}{\Gamma(1+a-c)} {}_1F_1(a; c; z) + \frac{\Gamma(c-1)}{\Gamma(a)} z^{1-c} {}_1F_1(1+a-c; 2-c; z),$$

as (see equations (2) and (4) on page 117 of [2])

$$D_\nu(z) = 2^{\frac{\nu}{2}} e^{-\frac{z^2}{4}} U\left(-\frac{\nu}{2}, \frac{1}{2}, \frac{z^2}{2}\right) = 2^{\frac{\nu-1}{2}} z e^{-\frac{z^2}{4}} U\left(\frac{1-\nu}{2}, \frac{3}{2}, \frac{z^2}{2}\right), \quad (5)$$

where in both cases  $-\pi < 2 \arg(z) \leq \pi$ . Applying the first formula in (5), we obtain

**Proposition 2.** *For the same parameter space as in the previous proposition, we have*

$$\begin{aligned} \mathcal{A}(p, q, r, s, t, u) &= 2^{-(p+1)} e^{i\pi u} (1 + e^{i\pi p}) \Gamma(p+1) \\ &\quad \times \lim_{a \downarrow 0} \mathbb{D}_a^{-u} \left[ \frac{e^{-at}}{(r+a)^{\frac{p+1}{2}}} U\left(\frac{p+1}{2}, \frac{1}{2}, \frac{(q-as)^2}{4(r+a)}\right) \right]. \end{aligned}$$

### 3. The integral $\mathcal{B}$

We apply again the Grünwald-Letnikov fractional derivative (4) of the exponential function. Firstly, we eliminate the denominator and the power  $x^p$  in the integrand, namely,

$$\mathcal{B}(p, q, r, s, t, u) = (-1)^u \mathbb{D}_s^u \mathbb{D}_{q-u}^p \left[ \int_{\mathbb{R}} \exp \left\{ (q-u)x - re^x - \frac{s e^x}{t + e^x} \right\} dx \right].$$

Now, we transform the inner integral  $I(s)$  and use the Maclaurin expansion with respect to  $s$  of the appropriate exponential term to obtain

$$\begin{aligned} & \mathcal{B}(p, q, r, s, t, u) \\ &= (-1)^u \mathbb{D}_s^u \left[ e^{-s} \mathbb{D}_{q-u}^p \left[ \int_{\mathbb{R}} \exp \left\{ (q-u)x - re^x + \frac{s t}{t + e^x} \right\} dx \right] \right] \\ &= (-1)^u \mathbb{D}_s^u \left[ e^{-s} \mathbb{D}_{q-u}^p \left[ \sum_{n \geq 0} \frac{(st)^n}{n!} \int_{\mathbb{R}} \frac{e^{(q-u)x - re^x}}{(t + e^x)^n} dx \right] \right] \\ &= (-1)^u \sum_{n \geq 0} \frac{t^n}{n!} \mathbb{D}_s^u \left[ e^{-s} s^n \mathbb{D}_{q-u}^p \left[ \int_{\mathbb{R}_+} \frac{y^{q-u-1} e^{-ry}}{(t+y)^n} dy \right] \right] \\ &= (-1)^u \sum_{n \geq 0} \frac{1}{n!} \mathbb{D}_s^u \left[ e^{-s} s^n \mathbb{D}_w^p [\Gamma(w) t^w U(w, w+1-n, r t)]_{w=q-u} \right], \quad (6) \end{aligned}$$

where in (6) the Laplace transform formula (see equation (2.1.3.1) on page 18 of [5])

$$\int_{\mathbb{R}_+} \frac{e^{-px} x^{\alpha-1}}{(x+z)^\rho} dx = \Gamma(\alpha) z^{\alpha-\rho} U(\alpha, \alpha+1-\rho, p z)$$

was used, which holds for all  $\Re(\alpha) > 0$ ,  $\Re(p) > 0$  and  $|\arg(z)| < \pi$ . This proves the following result.

**Proposition 3.** *For all  $p, s \geq 0, q \in \mathbb{R}, r, t, u > 0$  and  $q > -u$ , we have*

$$\begin{aligned} \mathcal{B}(p, q, r, s, t, u) &= e^{i\pi u} \sum_{n \geq 0} \frac{1}{n!} \mathbb{D}_s^u \left[ e^{-s} s^n \right. \\ &\quad \left. \times \mathbb{D}_{q-u}^p [\Gamma(q-u) t^{q-u} U(q-u, q-u+1-n, r t)] \right]. \quad (7) \end{aligned}$$

Unfortunately, our method holds true for  $q-u > 0$  only since  $\Gamma(q-u)$  in (7).

### 4. The integral $\mathcal{I}$

This time it is enough to split the integration domain into positive and negative reals and take the Maclaurin expansion in both sub-integrals of the exponential expression  $\exp\{-se^{-x}\}$ . We obtain the following

**Proposition 4.** For all  $p \geq 0, q \in \mathbb{R}, s, r > 0$ , we have

$$\mathcal{I}(p, q, r, s) = \frac{\Gamma(p+1)}{(2r)^{\frac{p+1}{2}}} \sum_{n \geq 0} \frac{(-s)^n}{n!} \left\{ D_{-p-1} \left( \frac{n-q}{\sqrt{2r}} \right) + e^{i\pi p} D_{-p-1} \left( \frac{q-n}{\sqrt{2r}} \right) \right\}. \quad (8)$$

Moreover, the following computable series representation holds:

$$\mathcal{I}(p, q, r, s) = \frac{\Gamma(p+1)}{(4r)^{\frac{p+1}{2}}} (1 + e^{i\pi p}) \sum_{n \geq 0} \frac{(-s)^n}{n!} U \left( \frac{p+1}{2}, \frac{1}{2}, \frac{(q-n)^2}{4r} \right).$$

Expression (8) could be deduced by some aspects of the discussion on pages 851-852 of [7]. However, there the authors' approach to quadratures for  $I(p, q, r, s)$  was completely different.

## Acknowledgement

The authors are indebted to anonymous referees for their careful reading of an earlier version of the manuscript and their valuable comments and corrections.

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