

MODELS WITH FEW DEGREES OF FREEDOM APPLIED TO SPECTRA OF LIGHT NUCLEI

N. Mankoč-Borštnik, M. V. Mihailović and M. Rosina

Faculty of Natural Sciences and Technology and J. Stefan Institute, University of Ljubljana, Ljubljana

In the study of nuclear structure one tries to find a small number of important degrees of freedom which are responsible for the features of low-lying spectra. One rather successful but expensive model for light nuclei is the complete diagonalization in a small model space¹⁾ where the effects of the rest of the space are approximately included in the effective interaction. The degrees of freedom refer to a small number of valence nucleons (typically 2-12) moving in a small single-particle space (typically 3 subshells). These valence nucleons are supposed to be responsible for the main features of several low-lying levels.

Our aim is to describe several levels with much fewer degrees of freedom. In the generator coordinate method (GCM)^{2,3)} we take for lower 1d2s nuclei two degrees of freedom - the quadrupole deformation β and the parameter μ . They are defined by means of the Nilsson-like single-particle field

$$\hat{h}(\beta, \mu) = \hbar\omega_0 \frac{1}{2}(-\nabla^2 + r) + \hbar\omega_0 \beta r^2 Y_{20} + C l s + \mu C l^2$$

which generates the single-particle states occupied in our Slater determinants $\phi(\beta, \mu)$. One then diagonalizes the Hamiltonian in the space of the wave functions

$$\Psi^J = \int \mathcal{P}^J(\beta, \mu) \mathcal{P}^J \phi(\beta, \mu) d\beta d\mu$$

Here \mathcal{P}^J is the projector on the subspace of good angular momentum J . The collective degrees of freedom " β -vibrations" and " μ -vibrations" are used to generate an appropriate N -body space of rather small dimension.

In the hermitian operator method (HOM)^{4,5)} we introduce as degrees of freedom all possible "particle-hole" excitations" with the respect to the ground state $|g\rangle$:

$$\sum_{ab} Q_{ab} (a_a^\dagger a_b + a_b^\dagger a_a) |g\rangle .$$

The Schrödinger equation in this subspace has the form

$$\sum_{cd} \left\{ \langle g | | a_a^\dagger a_b + a_b^\dagger a_a, | H, a_c^\dagger a_d + a_d^\dagger a_c | | | g \rangle \right\} Q_{cd} = \\ = \omega \sum_{cd} \left\{ \langle g | (a_a^\dagger a_b + a_b^\dagger a_a) (a_c^\dagger a_d + a_d^\dagger a_c) | g \rangle \right\} Q_{cd}$$

Here ω is the excitation energy and g should be a good approximation to the exact ground state in the given model space. The two-body density matrix of the ground state is needed as input information.

The generator coordinate method can be formally represented as describing a system of anharmonic, interacting phonons. In our case we have states with several phonons of two kinds (" β -phonons" and " μ -phonons"). On the other hand, the particle-hole states in the hermitian operator method can be formally represented as one-phonon states of several kinds of phonons (each level is a different phonon). Such representation may offer a bridge between our description of light nuclei and the description of heavier nuclei by the Alaga group⁶⁾.

In Fig. 1. some low-lying levels of ^{20}Ne are calculated with GCM²⁾ (even J , $T=0$) and HOM⁴⁾ (the lowest level for each J and T) and they are compared with the result of complete diagonalization¹⁾. For each J and T subspace of small dimension is taken (1-6 for GCM, 1-5 for HOM) compared to the complete diagonalization (3-56). We are preparing calculations of odd and even parity states in a larger single-particle space (3 major shells), using both GCM and HOM.

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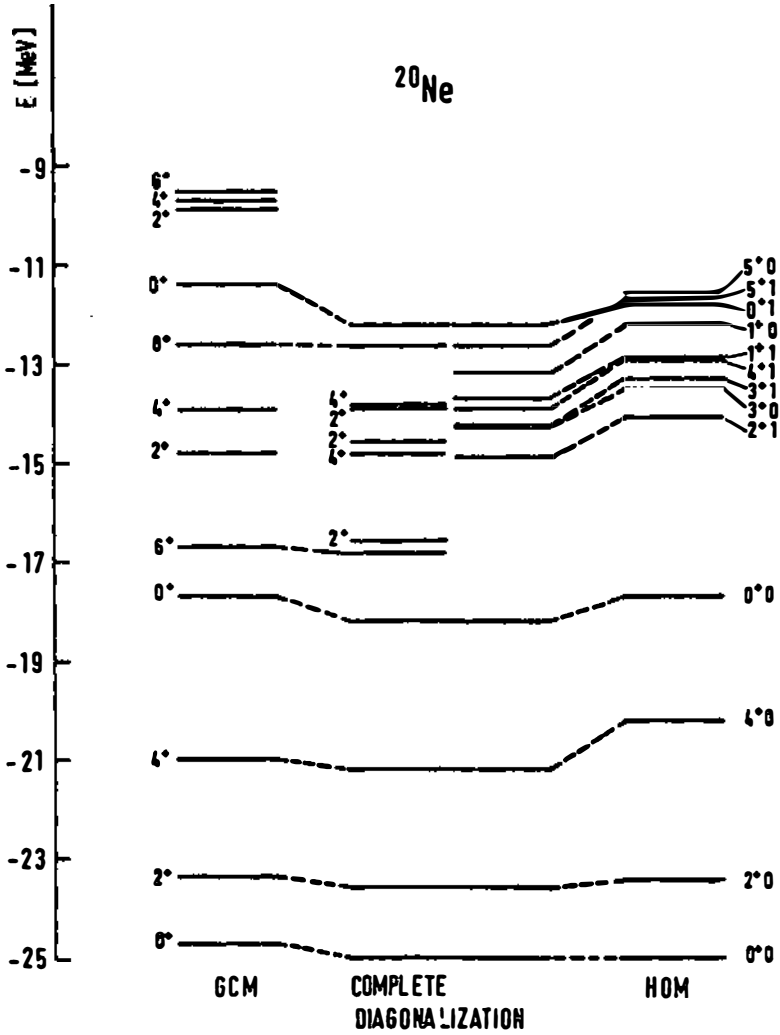


Fig. 1.