

E 16 High-Order Corrections to Electron Radial Wave Functions

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In recent years, radial wave functions for free¹⁾ and bound^{2,3)} electron states have been recalculated several times. One of the problems in such calculations is the choice of average potentials. Recently²⁾ capture ratios were calculated for different types of potentials and compared with experimental values. In this calculation the screening was represented by a linear combination of exponential functions, the parameters of which were determined by fitting the Hartree-Fock⁴⁾ (HF), Thomas-Fermi-Dirac⁵⁾ (TFD), or Herman-Shillman⁶⁾ (HS) potentials for a given atom.

1. Since the HS potentials are most extensively tabulated⁶⁾, we used them for the screening function and solved the Dirac one-electron radial equations. Extrapolation to the nuclear radius and the interpolation between the tabulated values of the HS potentials were performed. Our $(g_L/g_K)^2$ values for $Z > 30$ and the nuclear radius parameter $r_0 = 1.2$ agree within 1% with the values of Band and Zyryanova⁷⁾, as one can see from Table 1. This indicates that the

TABLE 1
 $(g_L/g_K)^2$

| <i>Z</i> | <i>BZ</i> | <i>Our</i> | <i>BB</i> | <i>Exp.</i> |
|----------|-----------|------------|-----------|-------------|
| 30 | 0.0987 | 0.0998 | 0.098 | 0.105 |
| 40 | 0.109 | 0.109 | 0.100 | 0.110 |
| 50 | 0.118 | 0.119 | 0.103 | 0.119 |
| 60 | 0.128 | 0.130 | 0.106 | 0.130 |
| 78 | 0.150 | 0.151 | 0.108 | 0.150 |

BZ ref. 7), $r_0 = 1.2$

Our our results, HS potential, $r_0 = 1.2$

BB from Fig. 1, ref. 2), HS potential, $r_0 = ?$

Exp from Fig. 2, ref. 2)

choice of the potential is not essential in this case. On the other hand, the HS potentials have been tabulated for all atoms, which is not the case with the HF and TFD potentials. It may happen that greater error is introduced by using a parameter fit to the relativistic HF and TFD potentials than by using an exact nonrelativistic HF, as in the HS tables.

2. We recalculated the ft values for $0^+ - 0^+$ transitions using free electron wave functions. A lot of work was done and some effects such as the compressibility⁸⁾ of the nucleus, were taken into account to obtain the uniformity of the ft values. (Table 2, column III).

Keeping the phase convention for β^+ decay correct, we obtain slightly different values (Table 2, column IV). Good agreement between the ft values is lost whether we take the correction to the nuclear matrix elements⁸⁾ (column V) into account, or not. Our calculations support the results of Bhalla⁹⁾ listed in column VI and show a greater difference between the ft values for ^{16}O and ^{34}Cl and those for other elements. These results will not affect the value of the vector coupling constant which is based on the decay of ^{16}O .

TABLE 2
 ft values for $0^+ \rightarrow 0^+$ transitions

| ft | Freeman | | Yamada | Our results | | Bhalla |
|------------------|-----------------------|------------------------|--------------------------|-------------|----------------|----------|
| | 10 ^{a)} I | 10 ^{b)} II | 8 ^{b,c)} III | a) IV | a-c) V | b) VI |
| ref. | | | | | | |
| ^{16}O | 3139 | 3127 | 3133 | 3138 | 3143 | 3127 |
| ^{26}Al | 3098 | 3086 | {3102 3131} | 3090 | {3102 3131} | 3085 |
| ^{34}Cl | 3150 | 3138 | 3141 | 3140 | 3140 | 3131 |
| ^{42}Sc | 3134 | 3122 | 3132 | 3111 | 3111 | 3109 |
| ^{46}V | 3143 | 3131 | 3133 | 3119 | 3113 | 3116 |
| ^{50}Mn | 3137 | 3125 | 3130 | 3104 | 3096 | — |
| ^{54}Co | 3144 | 3132 | 3131 | 3104 | 3094 | 3102 |

- a) Källen radiation corrections;
 b) Kinoshita-Sirkin radiation corrections;
 c) Nuclear matrix element corrections.

In all calculations, finite size and screening corrections are taken into account.

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