

The Inside-Triplet E2 Transition due to the Leading HRP  
 A Process for a Normal plus Superfluid System

V. Paar, Prirodoslovno-matematički fakultet, University of  
 Zagreb and "Rudjer Bošković" Institute, Zagreb, Yugoslavia

A. Ljubičić, "Rudjer Bošković" Institute, Zagreb, Yugo-  
 slavia

P.F. Bortignon, Istituto Galileo Galilei, Padova, Italy

We take into account second- and induced third-order diagrams for the E2 transition between two-phonon states, by combining the basic particle (quasiparticle)-vibration process in all possible ways. One vertex is always of the HRP type. This approach is exemplified for the  $4_1+2_2$  E2 transition in  $^{60}\text{Ni}$ . For protons, the system is normal, with proton particle-hole excitations appearing in the intrinsic phonon structure and for neutrons, it is a superfluid system with two quasineutron excitations appearing in the intrinsic phonon structure. The superfluid factors  $U_j, V_j$  and the quasiparticle energies for neutrons are the BCS solutions calculated using neutron single-particle energies of Reehel and Sorensen and the pairing strengths determined by fitting the odd-even mass difference. The  $B(E2)(4_1+2_2)$  is calculated as arising from the interference of the proton and quasi-neutron third-order HRP process. By applying the harmonic relation for  $B(E2)(4_1+2_1)$ , we calculated the intensity ratios

$$\frac{W(4_1+2_2)}{W(4_1+2_1)} = \frac{B(E2)(4_1+2_2)}{B(E2)(4_1+2_1)} \begin{pmatrix} 0.347 \\ 1.173 \end{pmatrix}^5$$

to be  $2.5 \times 10^{-5}$ ,  $5.8 \times 10^{-5}$  and  $11.9 \times 10^{-5}$  for the radial matrix elements in the physical region, with  $\langle r \frac{dv}{dr} \rangle$  equal to 40, 45 and 50 MeV, respectively. This is of the same order of magnitude as the experimental value  $(6.9 \pm 1) \times 10^{-5}$  (ref.1).

- 1) B.A. Logan, W.R. Dixon, R.S. Storey and A. Ljubičić,  
 Can. J. Phys. 55 (1977) 142.