by the presence of the higher symetry states, but for our purpose we can disregard D-state (see for example Ref. [4,5]) and charge form factor reads

$$F_{ch}(^3\mathrm{He}) = \left(F^p_{ch} + \frac{1}{2}F^n_{ch}\right)(P^2_s F_s + P^2_{s'} F_{s'}) - \frac{1}{2} P_s P_{s'} (F^p_{ch} - F^n_{ch})_{s-s'} \quad (1)$$

$$F_{ch}(^3\mathrm{H}) = (F^p_{ch} + 2F^n_{ch})(P^2_s F_s + P^2_{s'} F_{s'}) + P_s P_{s'} (F^p_{ch} - F^n_{ch}) F_{s-s'} \quad (2)$$

where the integrals F_s and $F_{s-s'}$ are given in Ref.[4] and in the Table, while $F_{s'}$ is given in Ref.[4] and is unimportante.

By direct inspection of the relations (1) and (2) and Table, one deduces the following conclusions:

1) for $P^2_{s'} = 0$, the diffraction minimum apears at $11(\mathrm{fm})^{-2} < q^2 < 12(\mathrm{fm})^{-2}$ for both ^{3}H and ^{3}He nuclei.

2) for $P^2_{s'} = 0.02$, the diffraction minimum is pushed back in q^2 for ^{3}He nucleus at $10(\mathrm{fm})^{-2} < q^2 < 11(\mathrm{fm})^{-2}$ and apears at $12(\mathrm{fm})^{-2} < q^2 < 13(\mathrm{fm})^{-2}$ for ^{3}H nucleus.

It is clear that the precise measurement of diffraction minimum for ^{3}H nucleus will determine the $P^2_{s'}$.

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4.7. Preliminary report about the work on ^{6}He + ^{1}H

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The cross-section of ^{1}H(^{6}He, t)^{4}He reaction is determined for $E_{^6\mathrm{He}} = 0 - 5$ MeV by irradiating CH and CH_2 targets with ^{6}He particles obtained from the primary reaction ^{7}Li(t, ^{6}He)^{4}He. The relatively large differential cross section of 0.7 b/sr is explained by low Coulomb barrier and a probable two-neutron correlation.

4.8. ^{6}Li(t, p)^{8}Li reaction at low energies*

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